

Complex stochastic dynamics in spin glasses and optimization problems

Federico RICCI TERSENGHI

Dip. Fisica, Univ. La Sapienza, Roma

in collaboration (over the years) with:

Giorgio PARISI, Enzo MARINARI,

Riccardo ZECCHINA, Silvio FRANZ,

Marc MEZARD, Andrea MONTANARI,

Guilhem SEMERJIAN, Florent KRZAKALA,

Lenka ZDEBOROVA, Tommaso RIZZO,

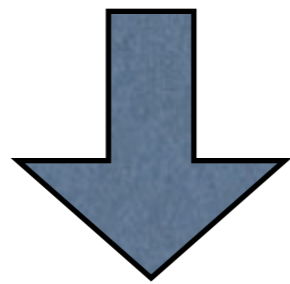
Spin Glasses

$$\mathcal{H} = - \sum_{\langle ij \rangle} J_{ij} s_i s_j$$

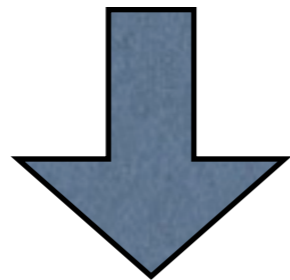
- $s_i = \pm 1$ classical (Ising) spins
no quantum effects!
- J_{ij} quenched random variables
e.g. $J_{ij} = \pm 1$ or $J_{ij} \sim N(0, 1)$

Spin Glasses

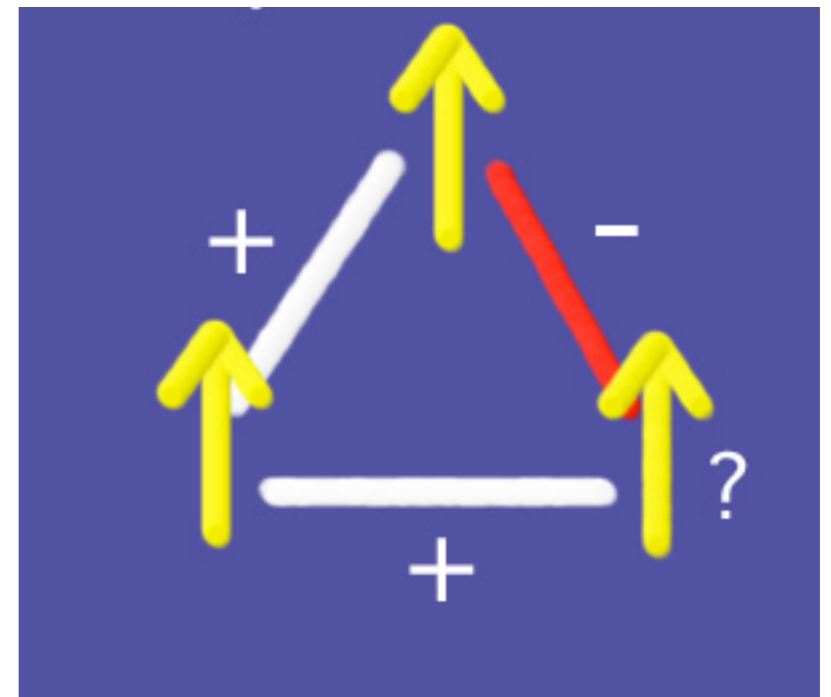
DISORDER



FRUSTRATION

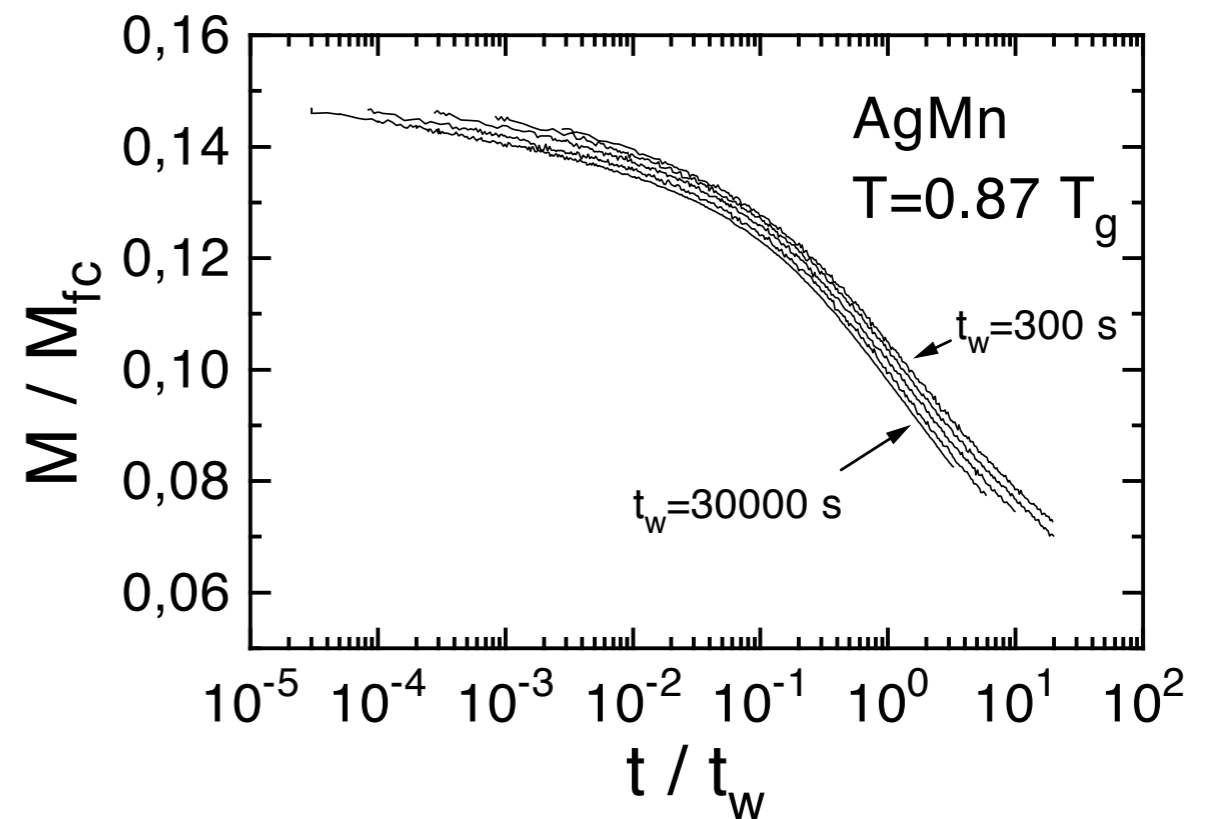
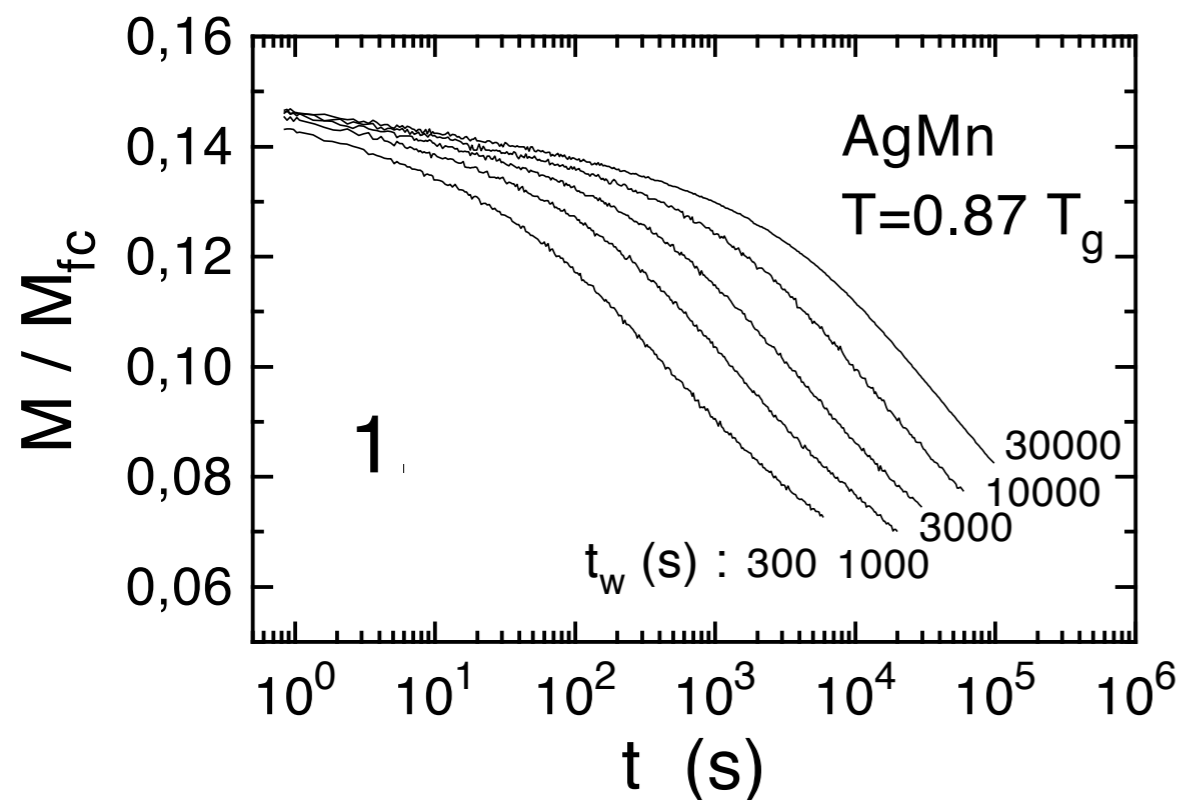


SLOW DYNAMICS



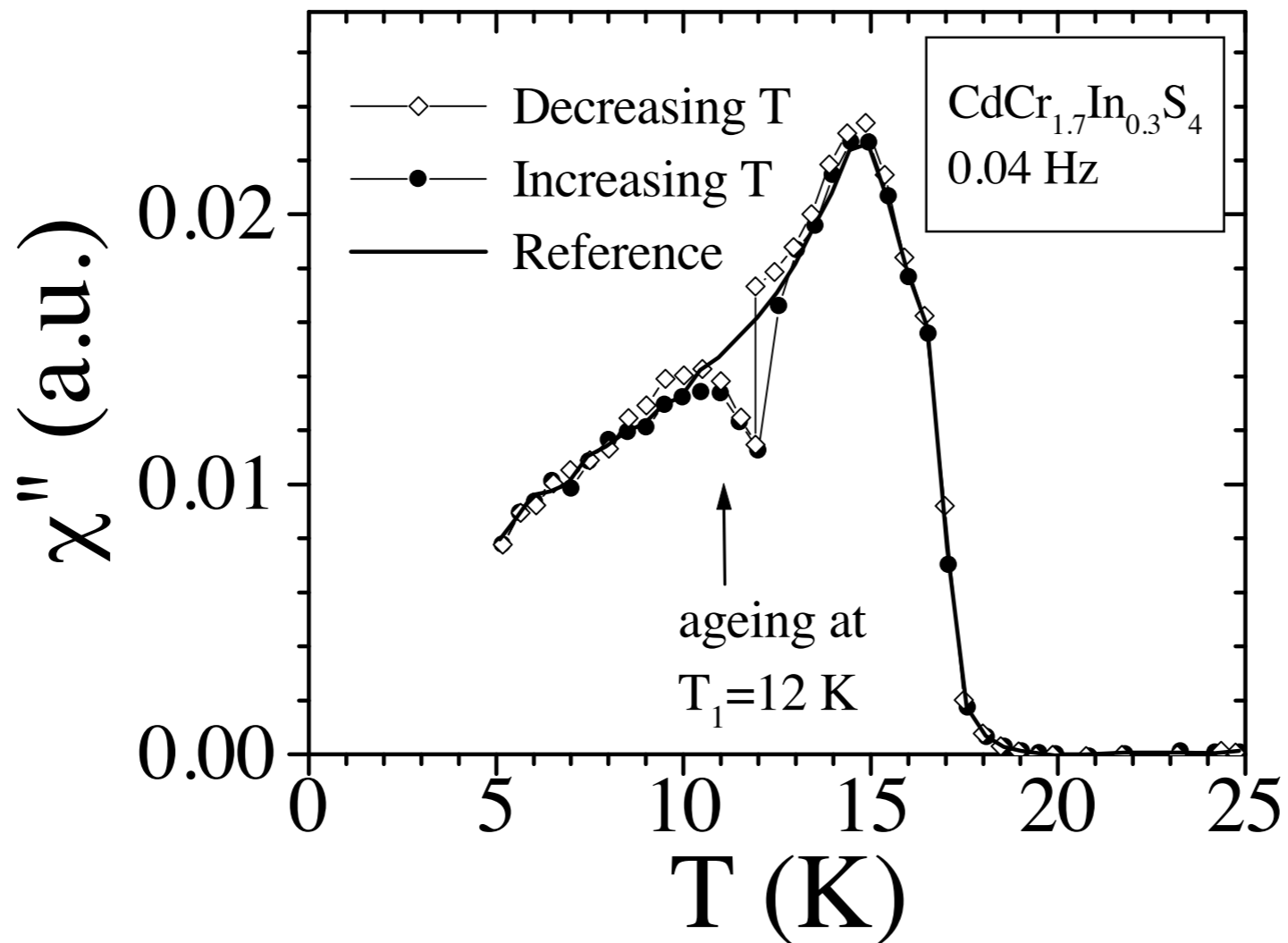
Experiments on SG

Below T_c \rightarrow always out of equilibrium



Experiments on SG

Rejuvenation and memory



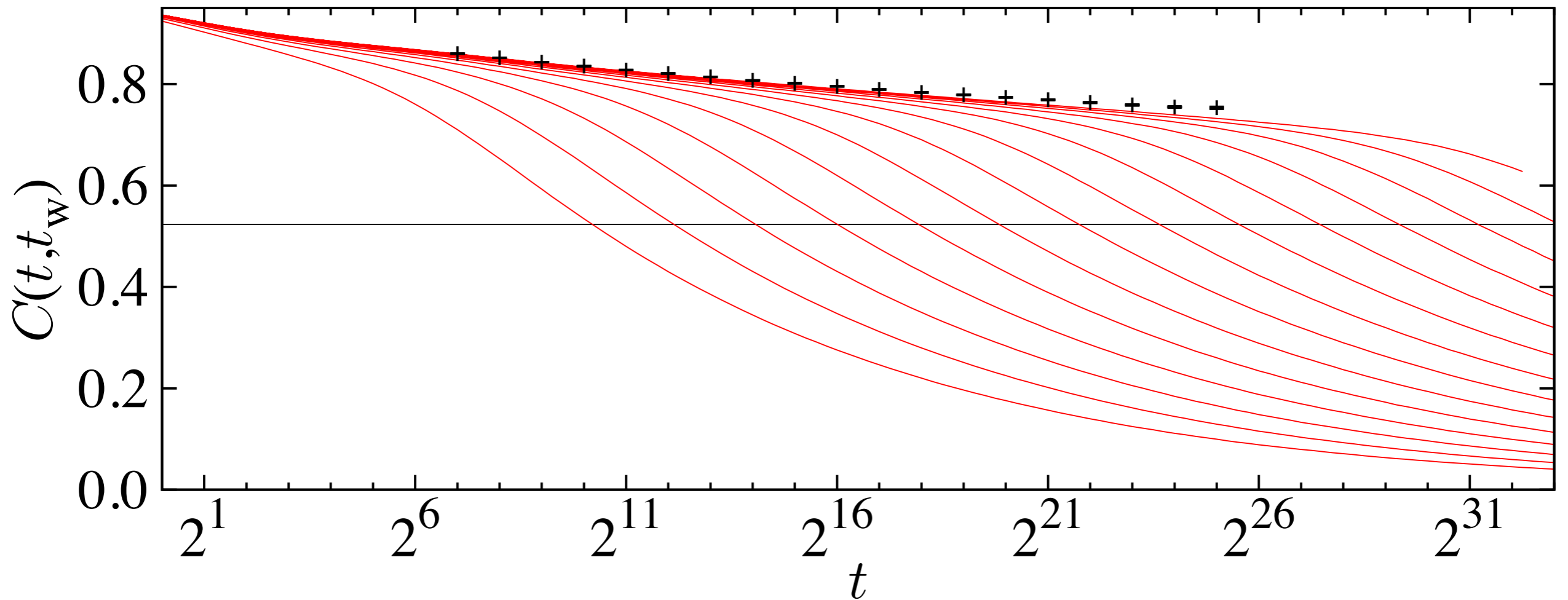
Slow growth of different kind of orders
at different temperatures

Complex macroscopic dynamics but simple microscopic rules

- Energy relaxation
- Local in the configurational space
- Stochastic in nature

Monte Carlo, Glauber, Langevin,

Numerical simulations of SG



$$C(t, t_w) = \frac{1}{N} \sum_i s_i(t_w) s_i(t_w + t)$$

3D SG $J=+/-1$
 $T=0.64T_c$

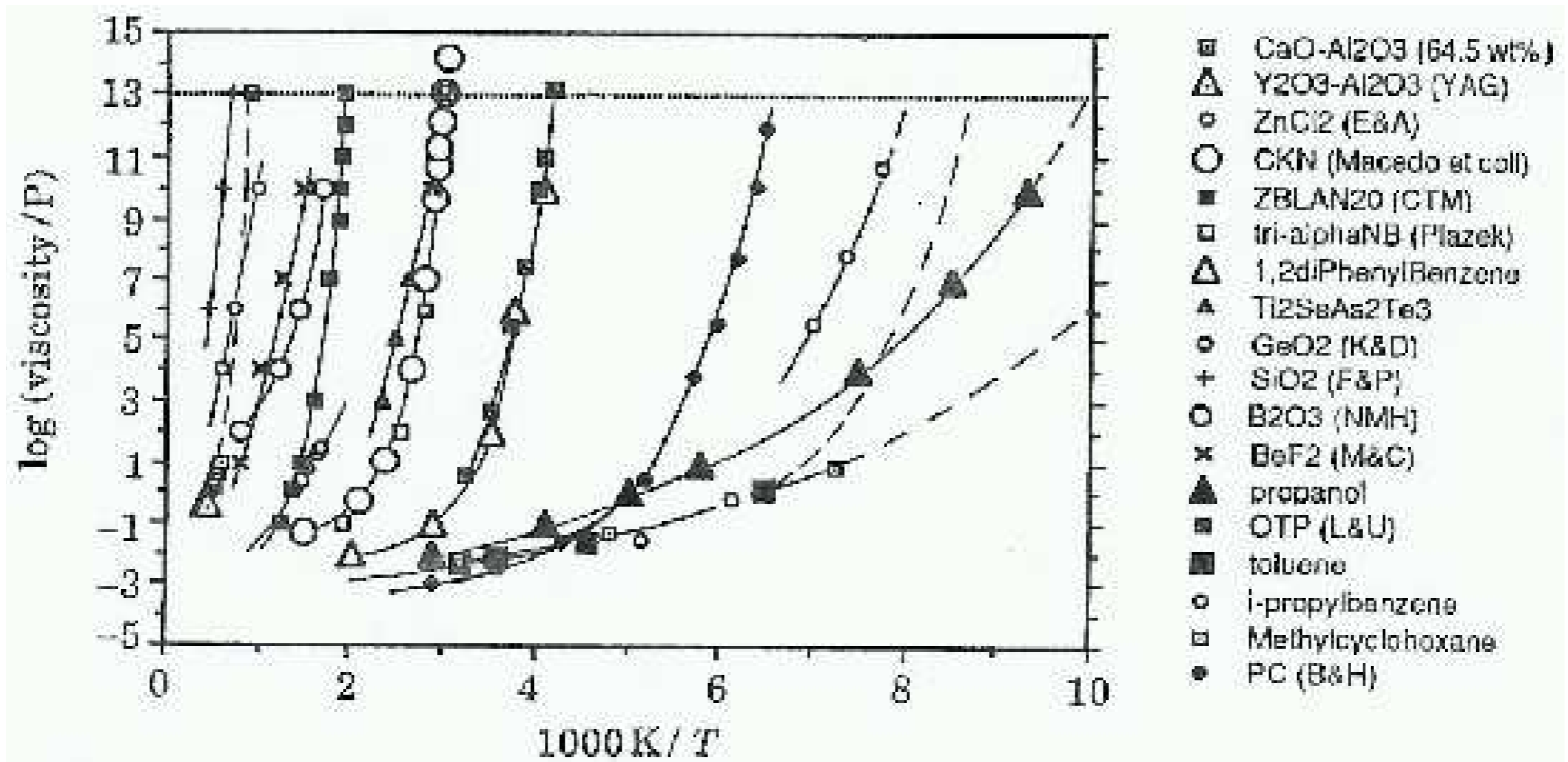
p-Spin Glass Model

$$\mathcal{H} = - \sum_{\langle ijk \rangle} J_{ijk} \underbrace{S_i S_j S_k}_{p\text{-uples}}$$

$p > 2$ is very different from $p = 2$

Random First Order Transition (RFOT)

(Structural) Glasses



Impressive increase of relaxation timescales!

A Solvable pSG Model

- Fully connected topology
- Continuous unbounded variables $S_i \in \mathbb{R}$
- Spherical constraint $\sum_{i=1}^N S_i^2 = N$
- Langevin dynamics

$$\partial_t S_i(t) = -\frac{\partial \mathcal{H}}{\partial S_i} - \mu(t) S_i(t) + \eta_i(t)$$

A Solvable pSG Model

- Closed integral differential eqs. for

$$C(t, t') = \frac{1}{N} \sum_i \langle S_i(t) S_i(t') \rangle \quad R(t, t') = \frac{1}{N} \sum_i \frac{\partial \langle S_i(t) \rangle}{\partial H(t')}$$

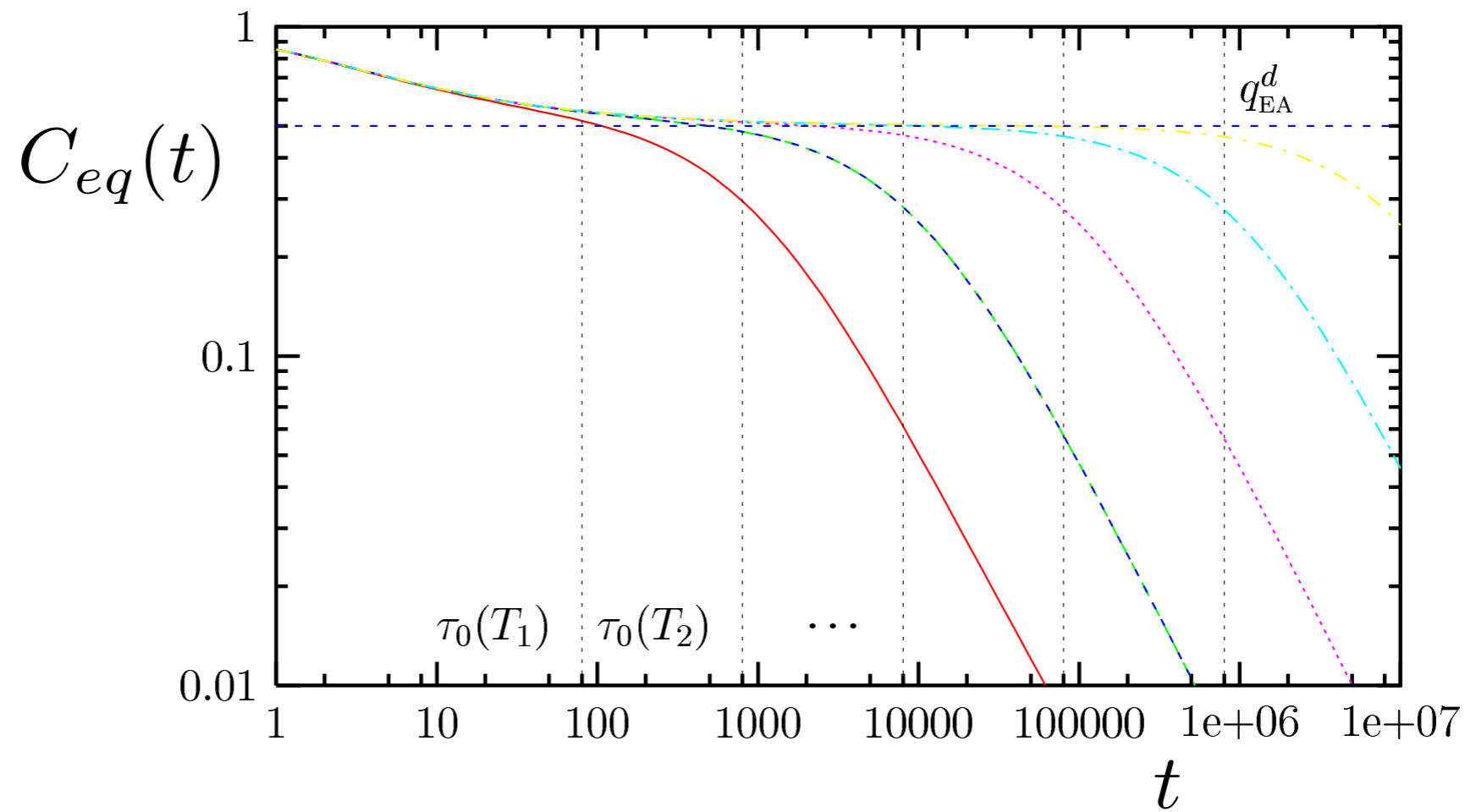
- For $T > T_c$ reaches equilibrium

$$C(t, t') = C_{eq}(t - t')$$

- For $T < T_c$ out of equilibrium (aging)

$$C(t, t') = \tilde{C}(h(t)/h(t')) \sim \tilde{C}(t/t')$$

pSG Dynamics for $T \rightarrow T_c^+$

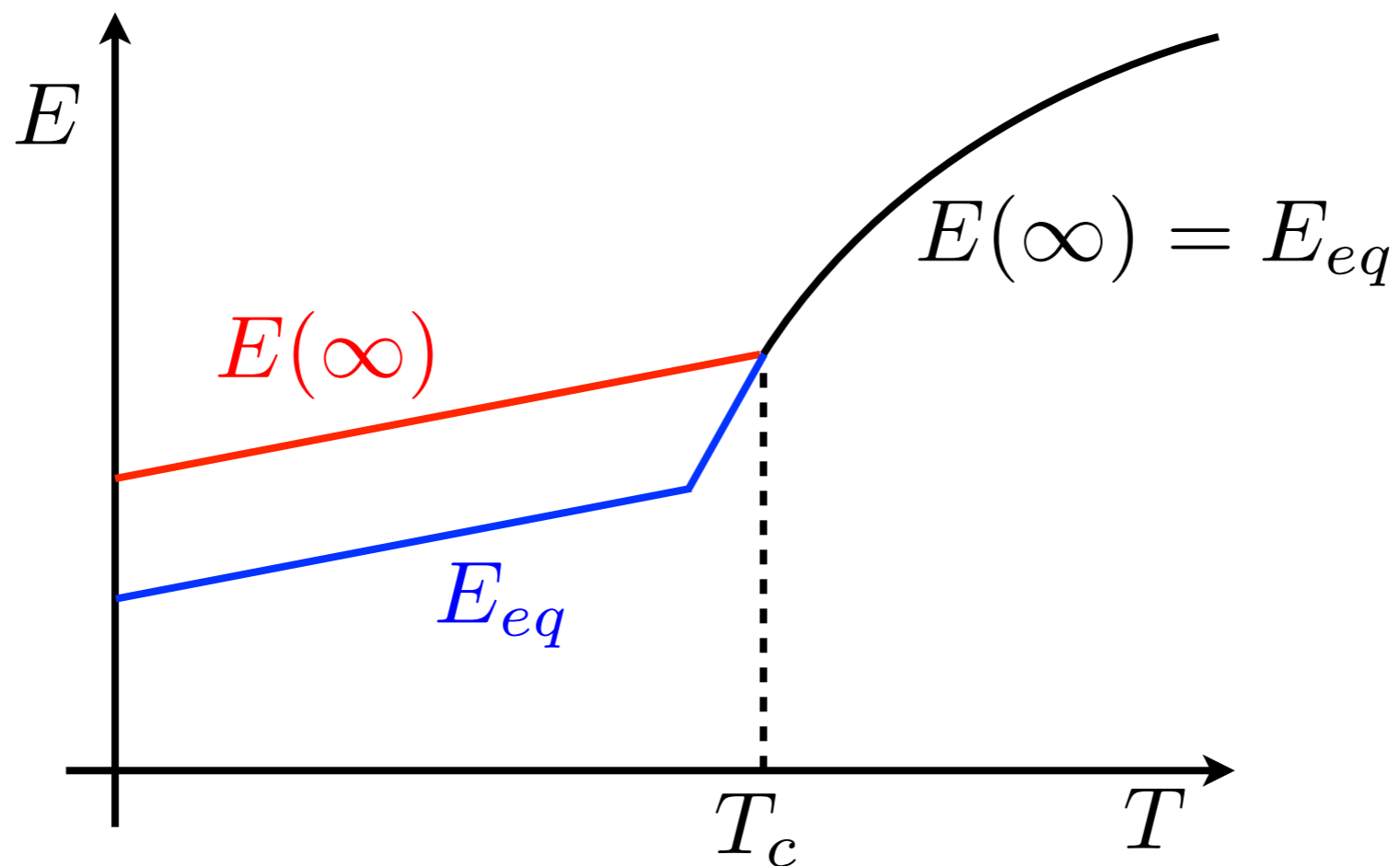


$$\tau(T) \propto (T - T_c)^{-\gamma}$$

pSG Dynamics for $T < T_c$

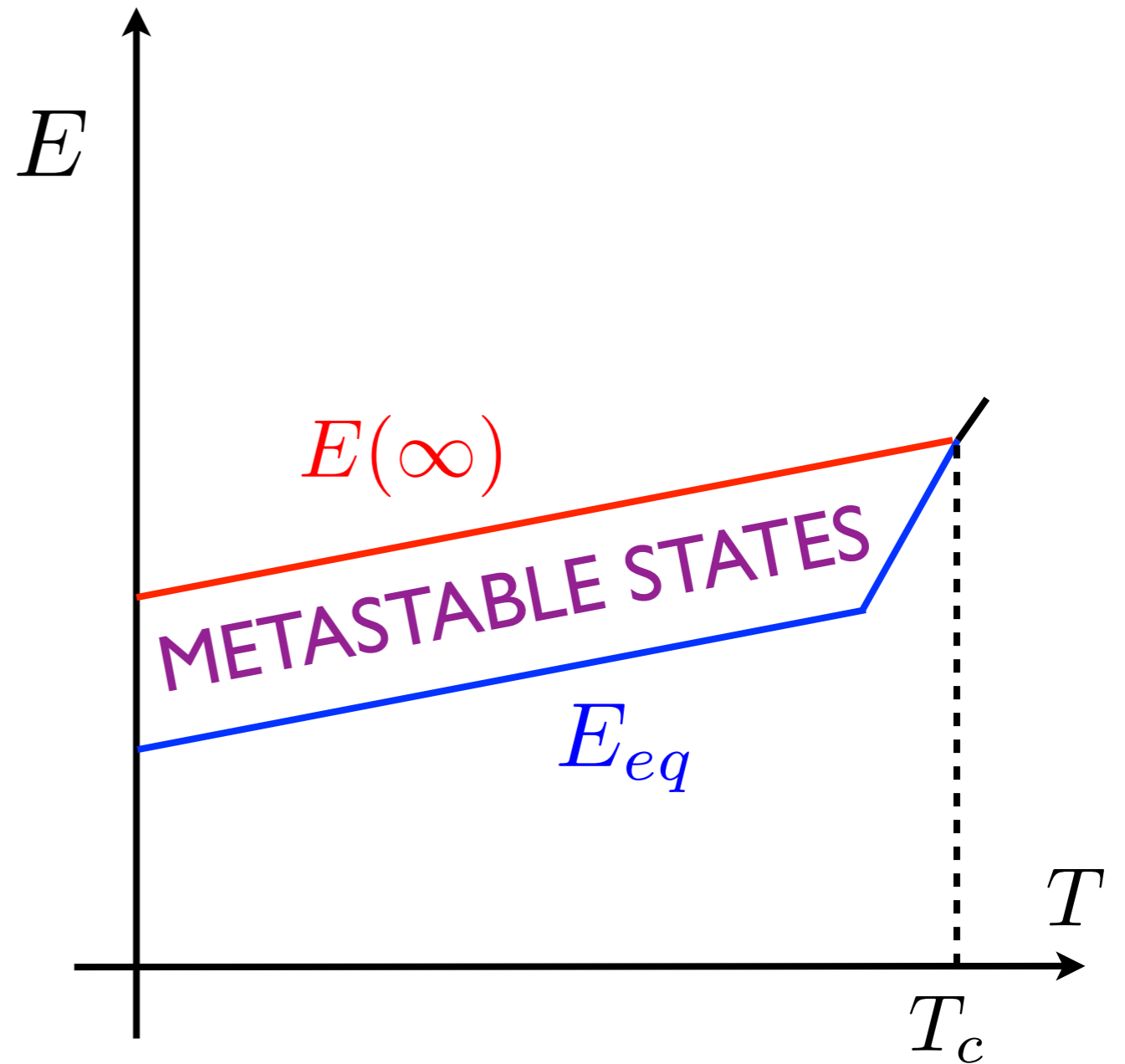
Does not relax to equilibrium states!!

$$E(\infty) \equiv \lim_{t \rightarrow \infty} E(t) \neq E_{eq}$$

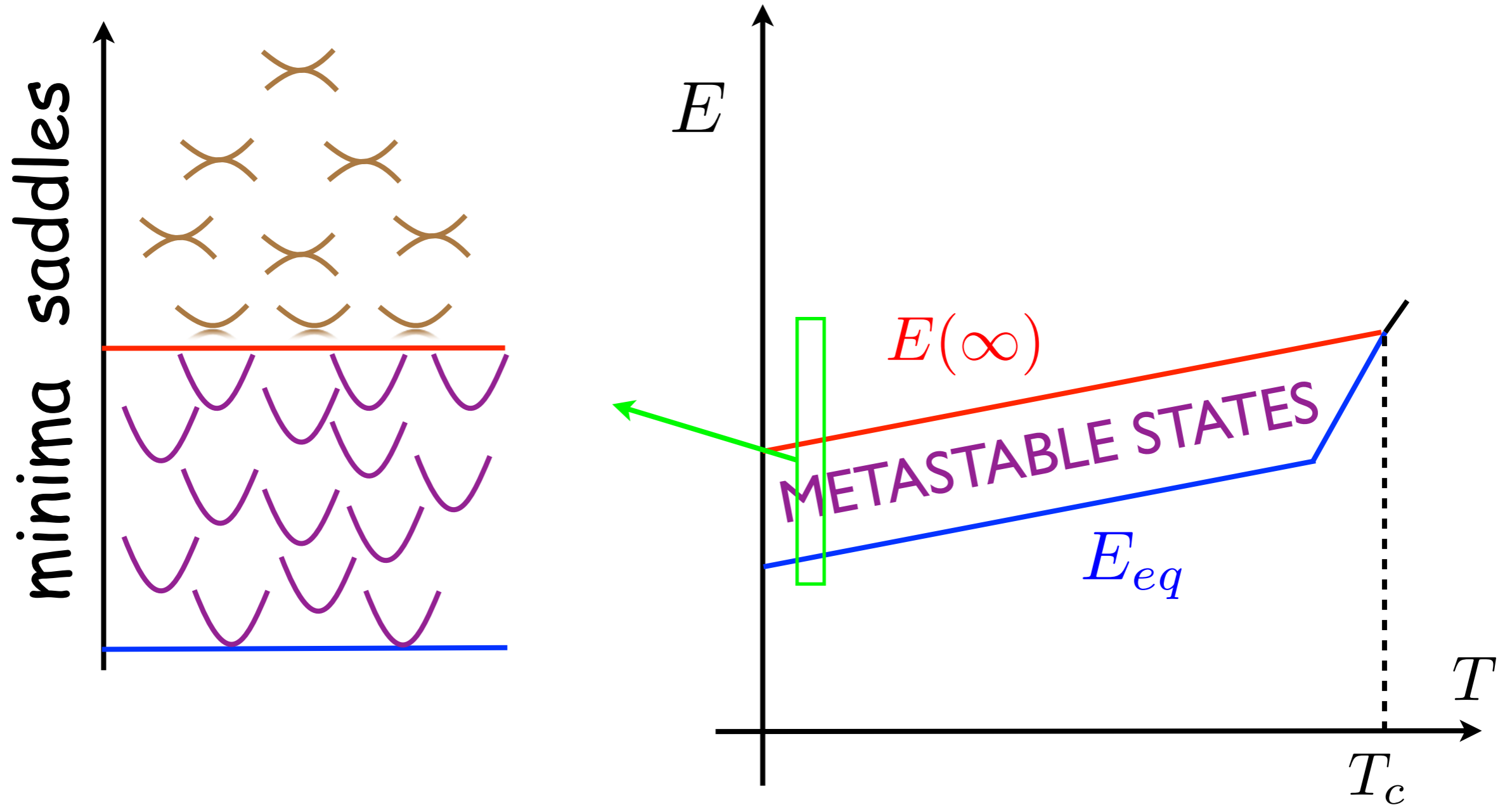


both for
quenches
& coolings

Why dynamics get stuck?



Why dynamics get stuck?



Static-Dynamics Connection ?

- It is too difficult to solve in general the dynamics (e.g. for Ising spins)
- Is long time dynamics determined by few thermodynamical (i.e. static) properties of the energy potential?
- E.g. does energy relaxation stop at the highest metastable states?

How to count metastable states

Compute the replicated free-energy $\Phi(m, T)$
by the replica or the cavity method

$$e^{-\beta m \Phi(m, T) N} \equiv \sum_{\alpha} Z_{\alpha}^m = \int e^{-\beta m f N + N \Sigma_f(f, T)} \mathrm{d}f$$

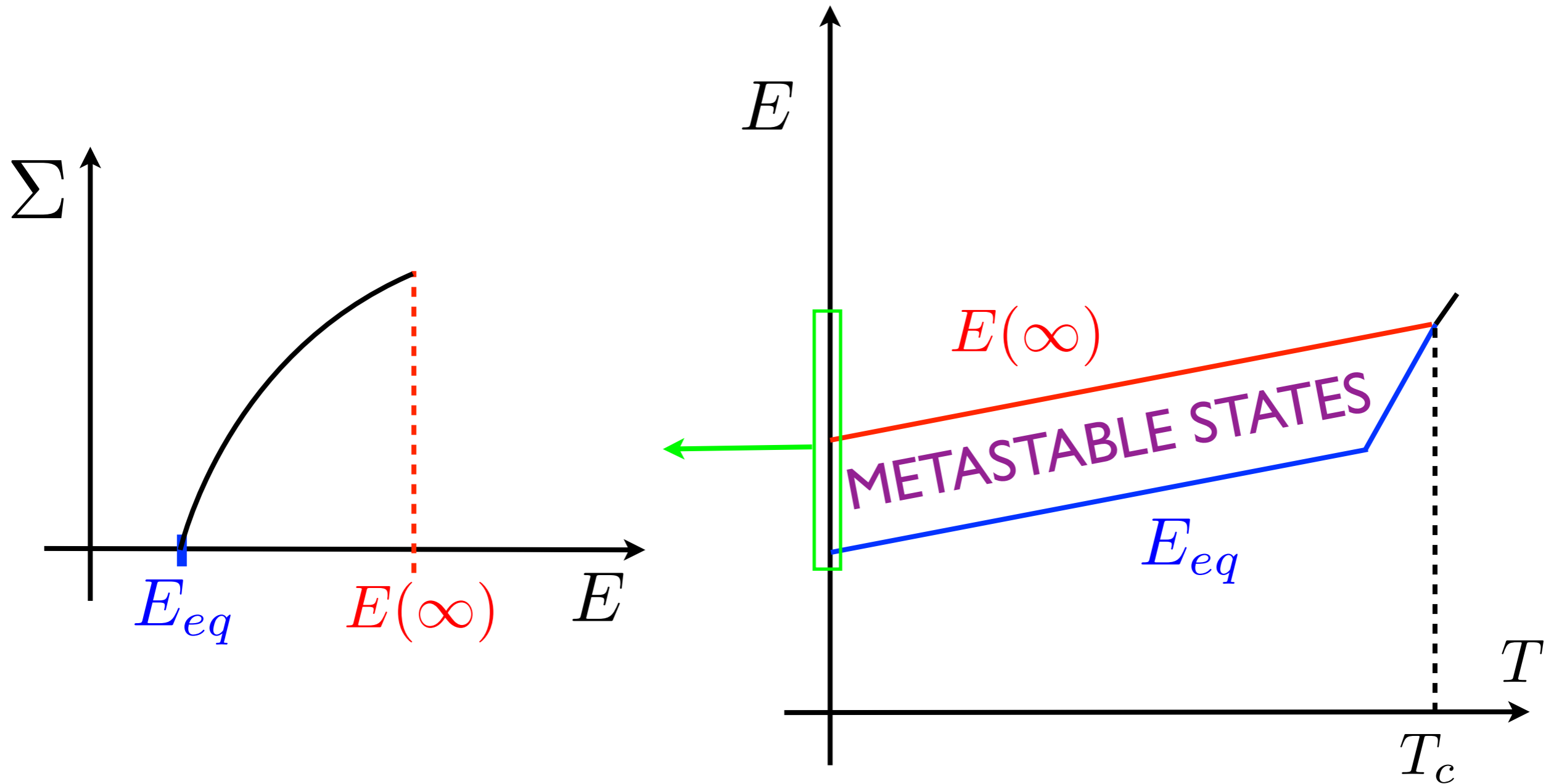
get the complexity or configurational entropy

$$\Sigma_f(f, T) = \ln(\# \text{ states with free-energy } f)$$

by Legendre transforming

$$\Sigma_f(f, T) = \beta m f - \beta m \Phi(m, T) \Big|_{\beta f = \partial_m (m \Phi)}$$

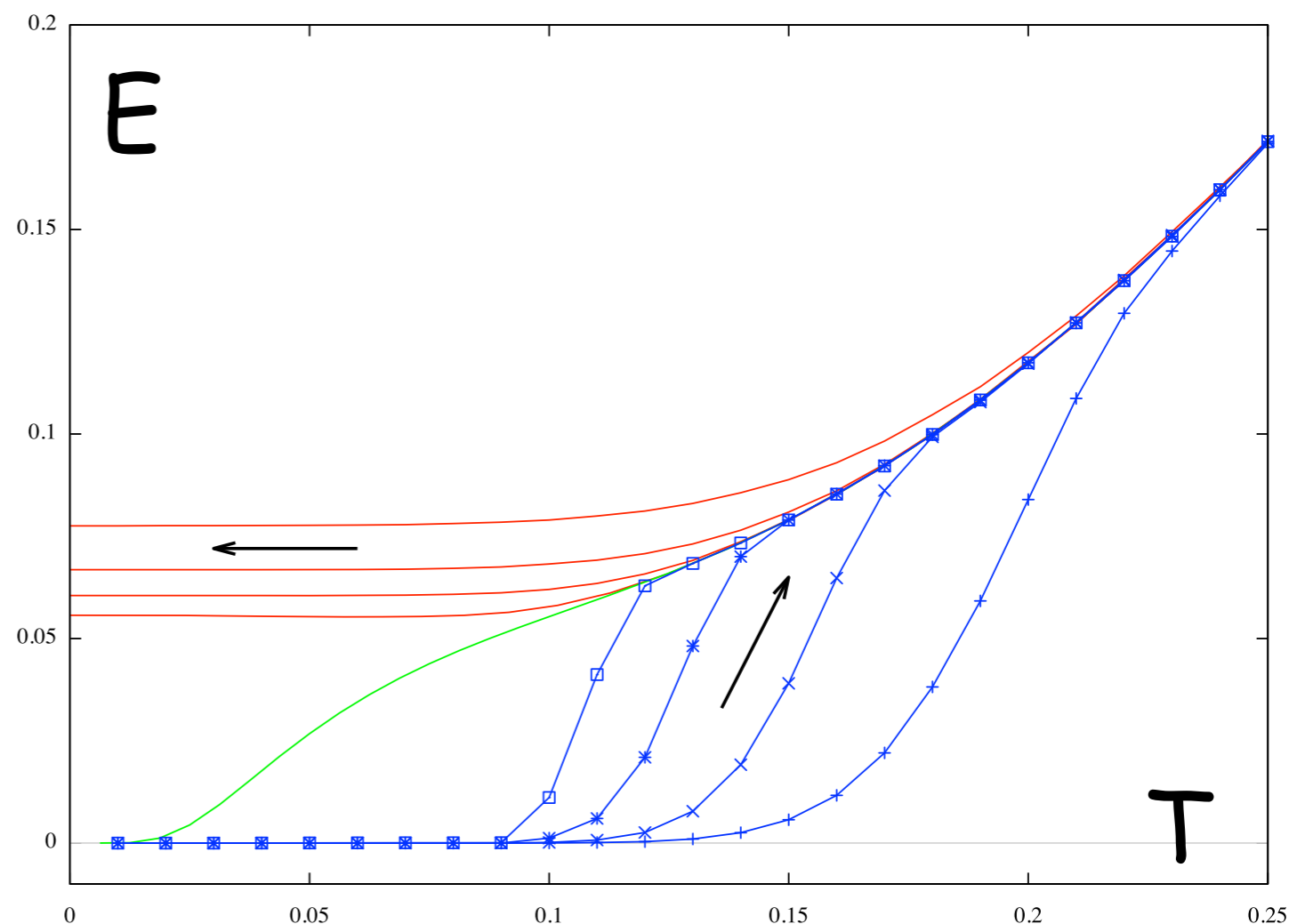
Why dynamics get stuck?



...and beyond mean field ?

Metastable states have infinite lifetime in MF
...but even in finite dimensions they can be huge!

1D model with local
4-spin interactions
(no disorder at all)



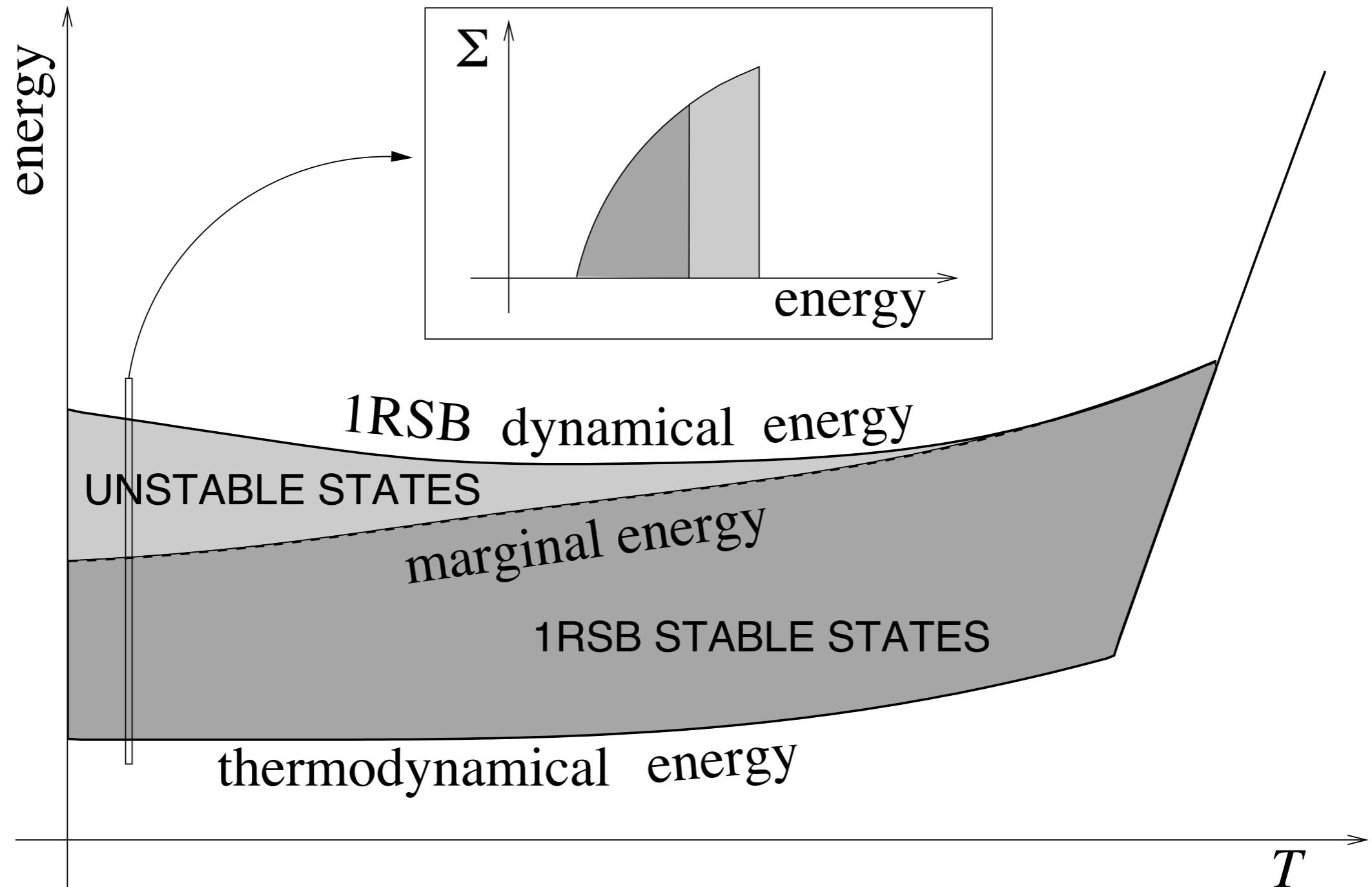
Ising SG are more complex

$$\mathcal{H} = - \sum_{\langle ijk \rangle} J_{ijk} S_i S_j S_k \quad S_i = \pm 1$$

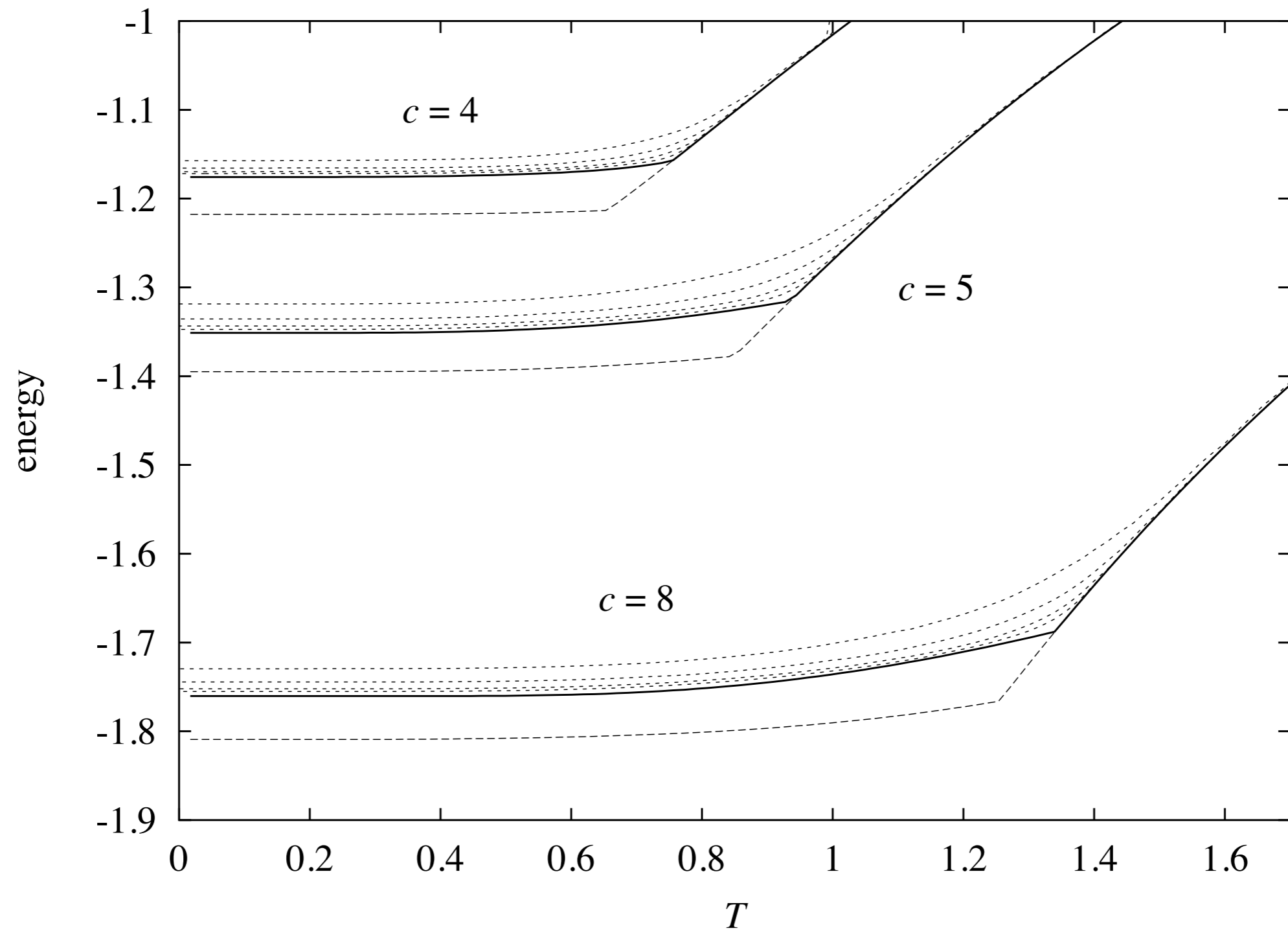
$$J = +/-1 \quad p=3$$

interaction network \rightarrow random c -regular graph:
 $\langle ijk \rangle$ randomly chosen such that each spin
participate exactly to c interactions

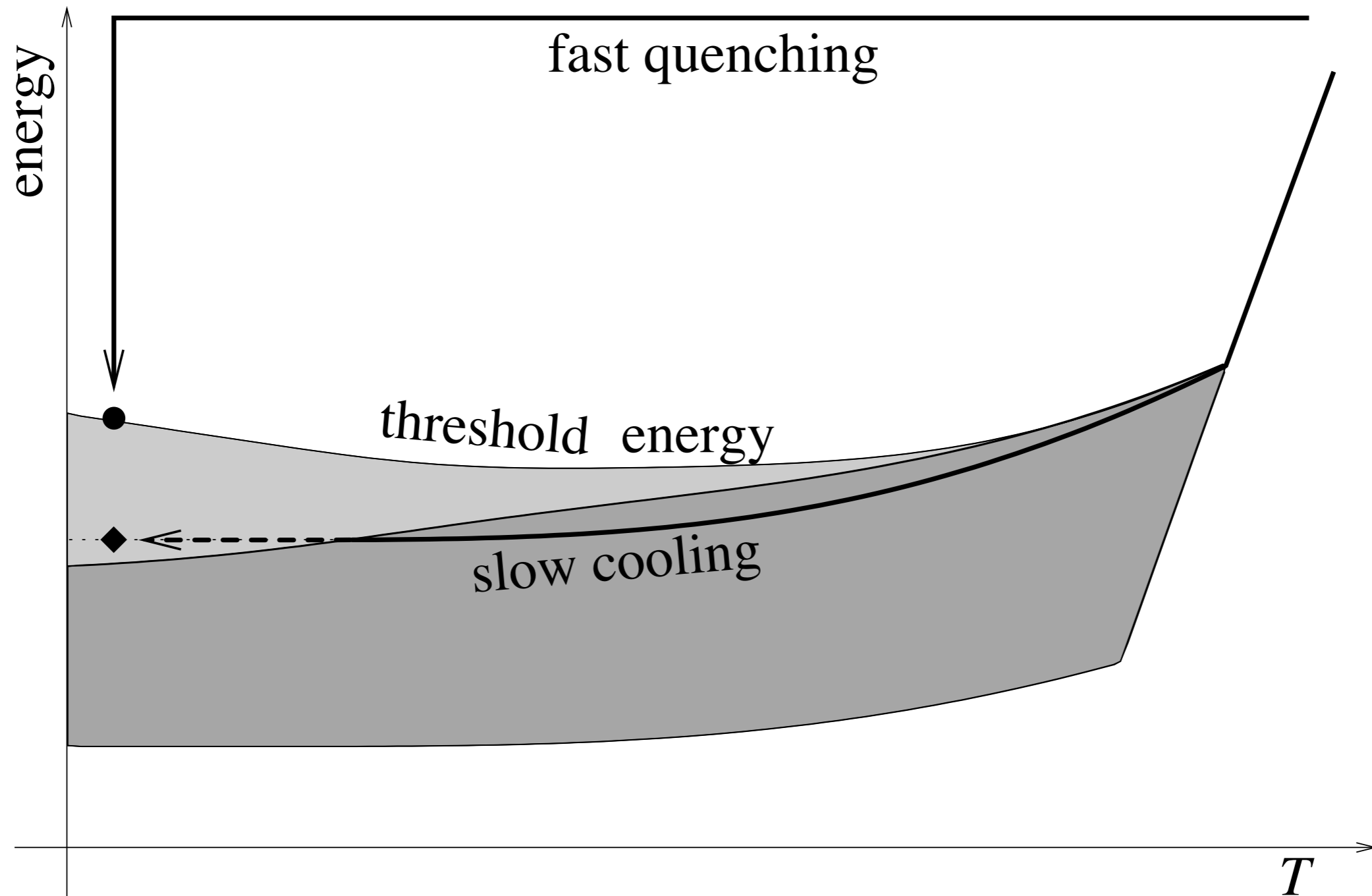
More Complex Complexity >:-!



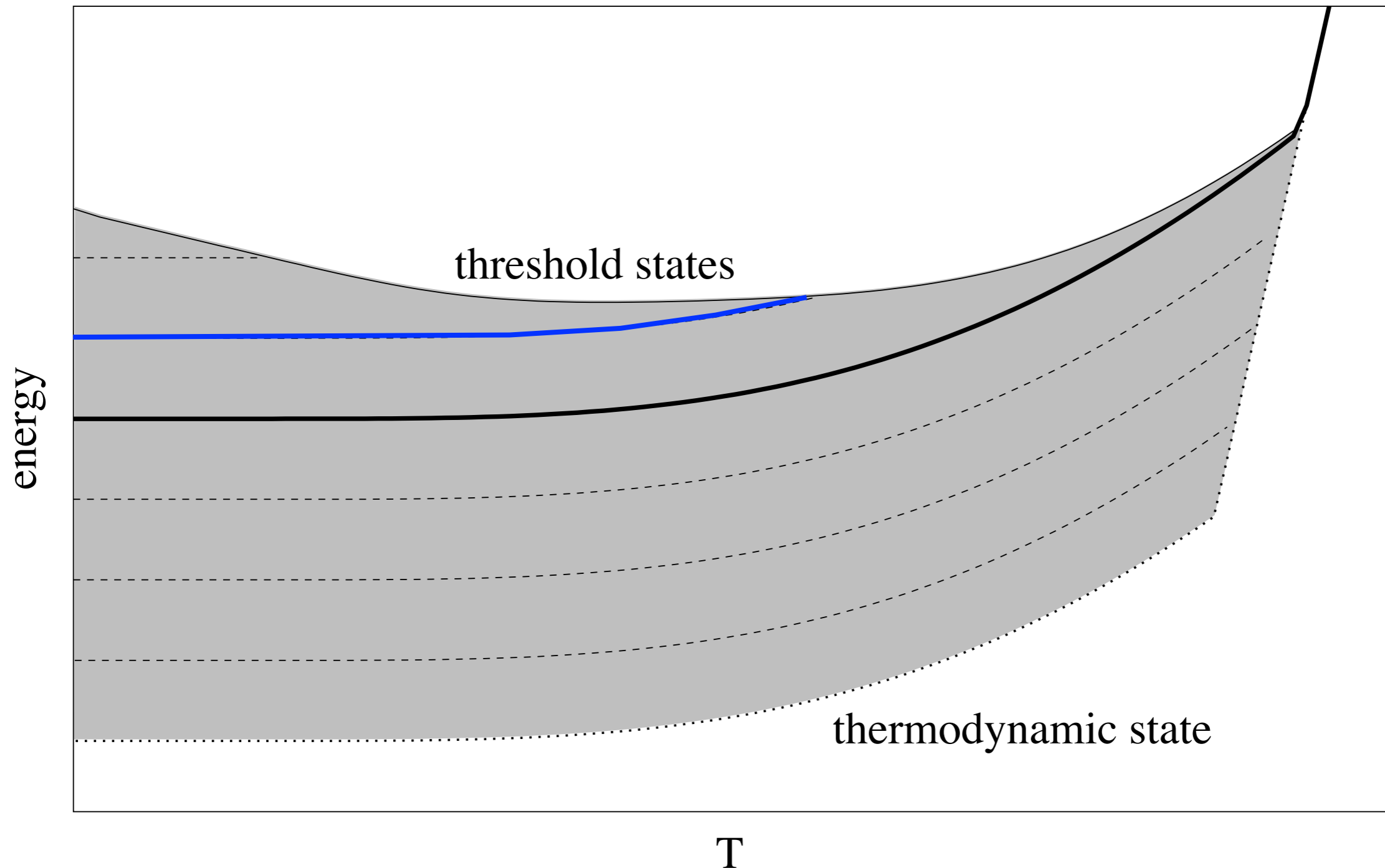
Coolings get stuck, as usual...



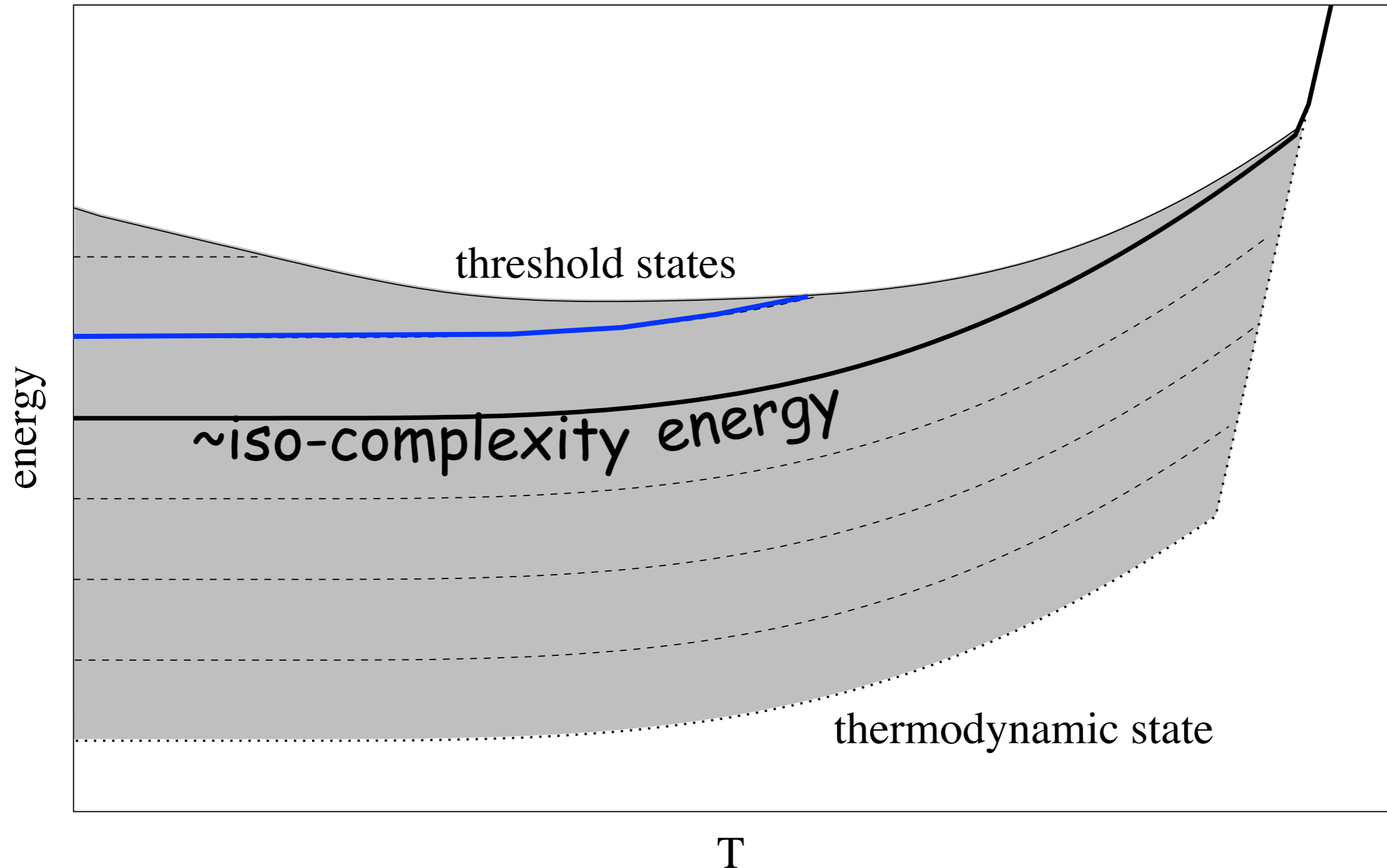
Coolings vs. Quenches



Coolings follow the states...



Coolings follow the states...



Another Static-Dynamics Link

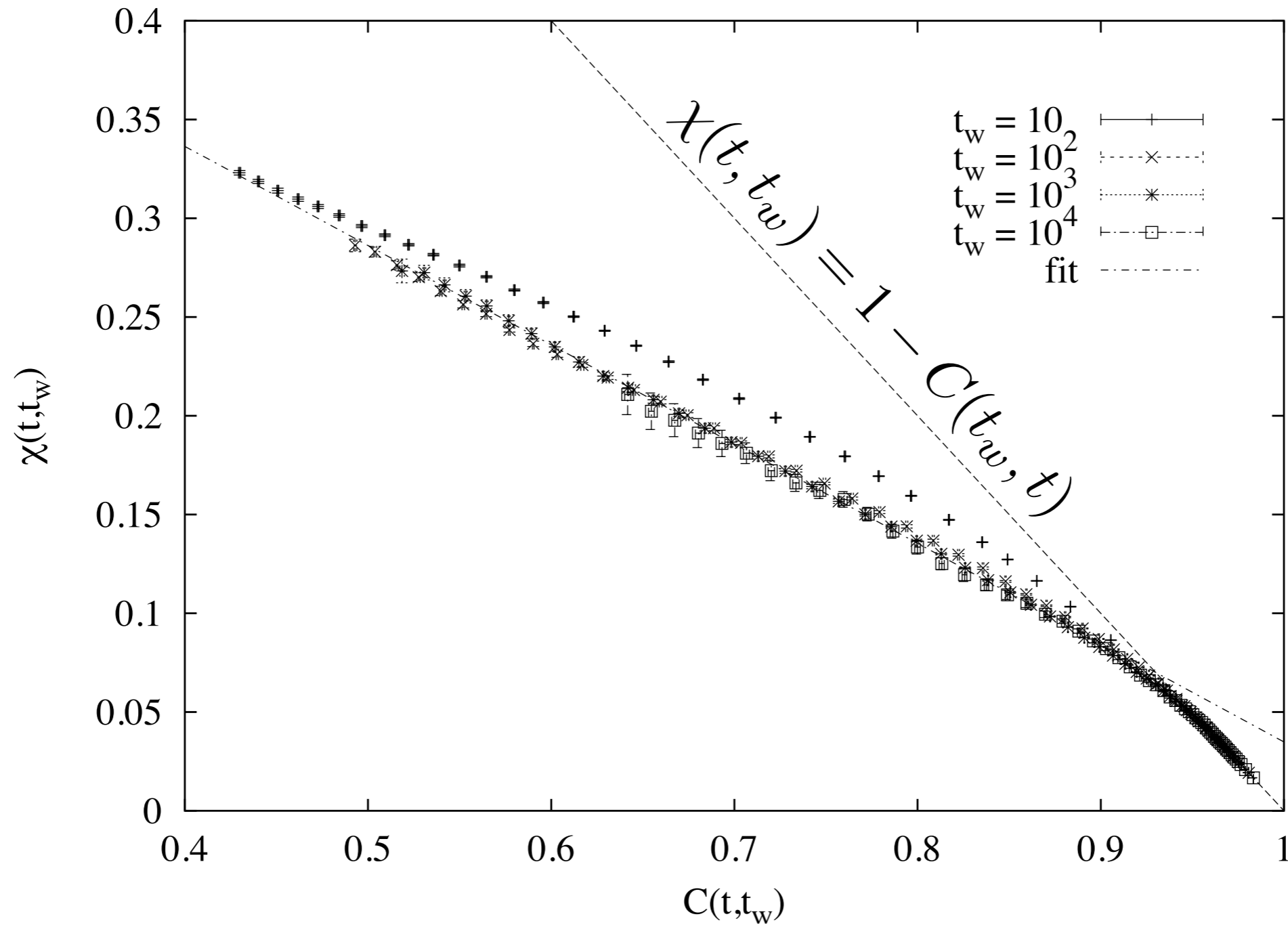
Measure correlation and integrated response

$$C(t_w, t), \quad \chi(t, t_w) = T \int_{t_w}^t R(t, s) ds$$

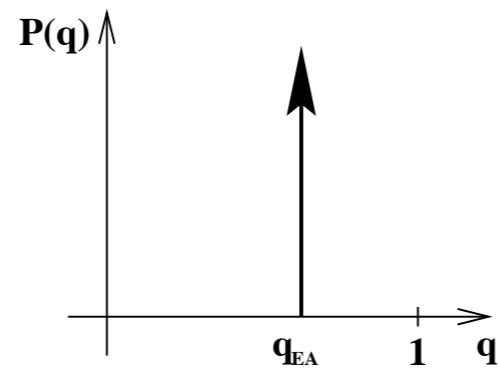
At equilibrium Fluctuation Dissipation Th.

$$\chi(t, t_w) = 1 - C(t_w, t)$$

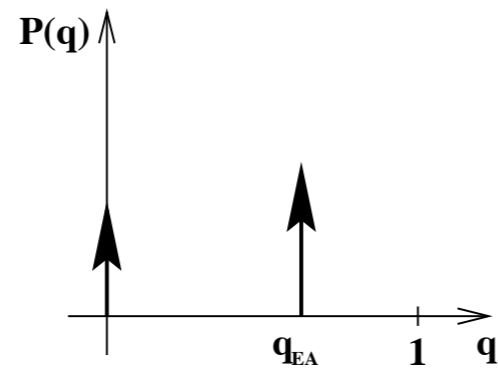
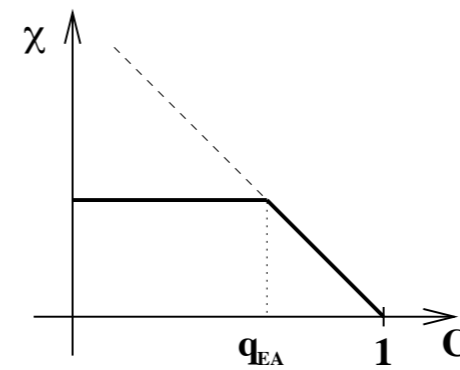
Fluctuation Dissipation Ratio



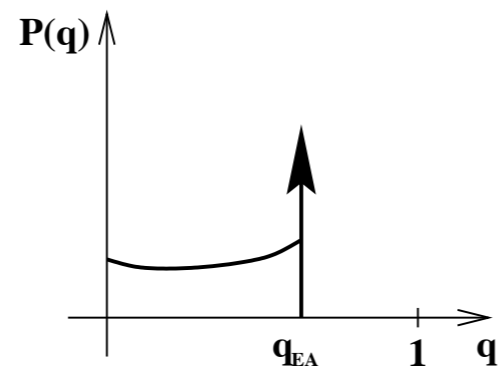
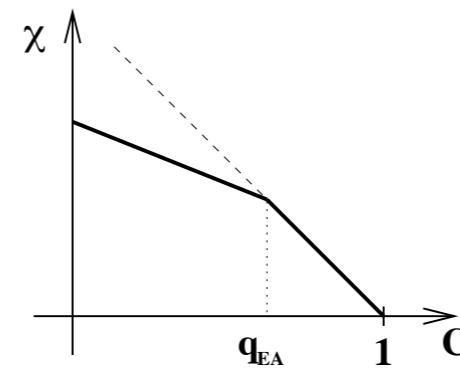
SG order parameter vs. FDR



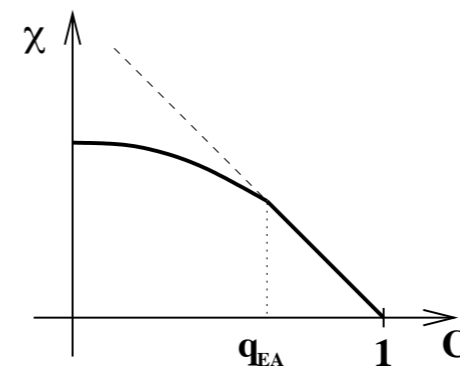
A



B



C



Combinatorial Optimization

Constraint Satisfaction Problems (CSP)

Find an assignment to N binary variables
such as to satisfy M constraints.

(each constraint involves few variables)

$$\alpha = \frac{M}{N}$$

constraints per
variable ratio

SAT/UNSAT threshold $\left\{ \begin{array}{ll} \alpha < \alpha_s & \text{SAT} \\ \alpha > \alpha_s & \text{UNSAT} \end{array} \right.$

Random 3-xorsat

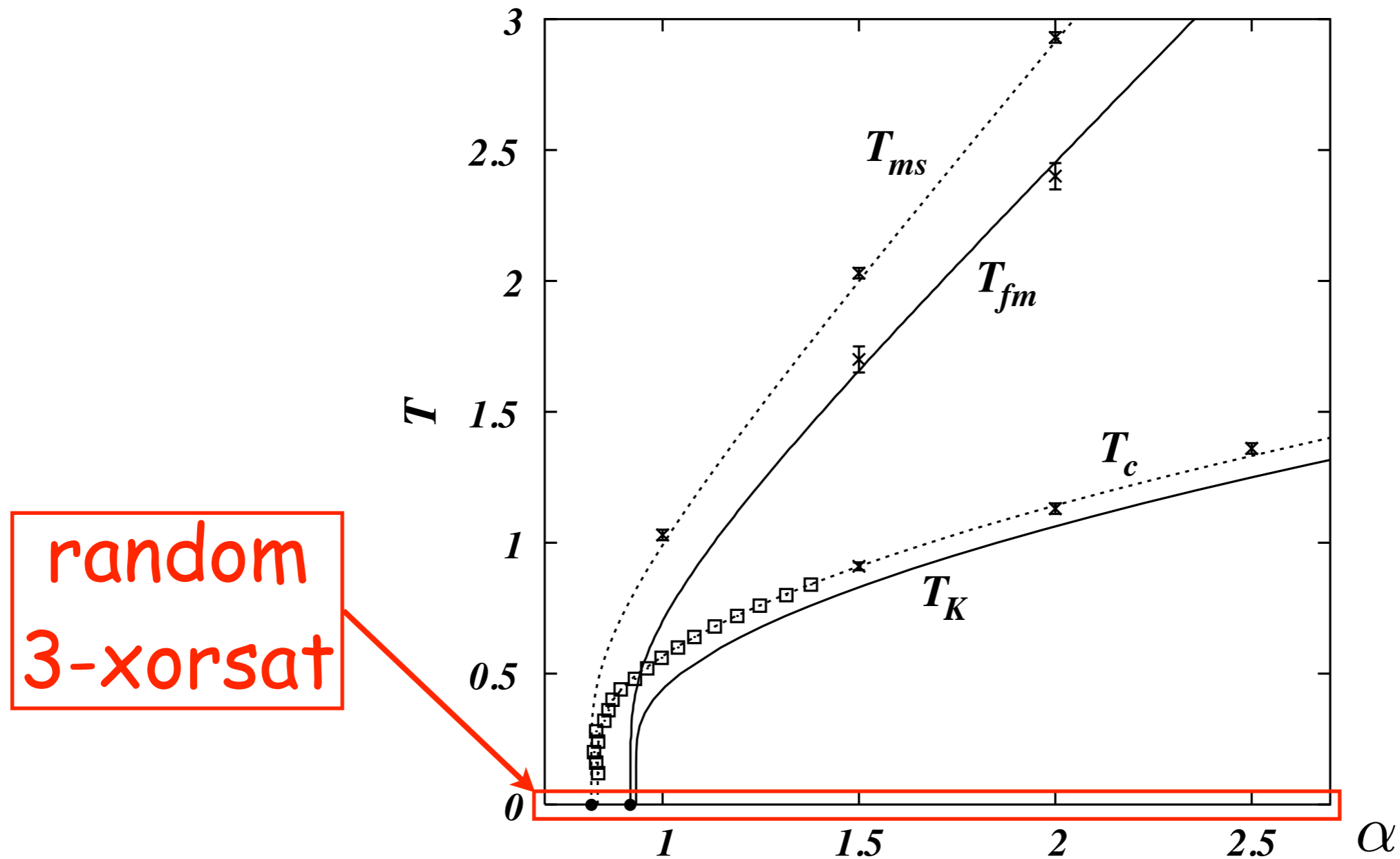
Satisfy αN equations like

$$S_i S_j S_k = J_{ijk}$$

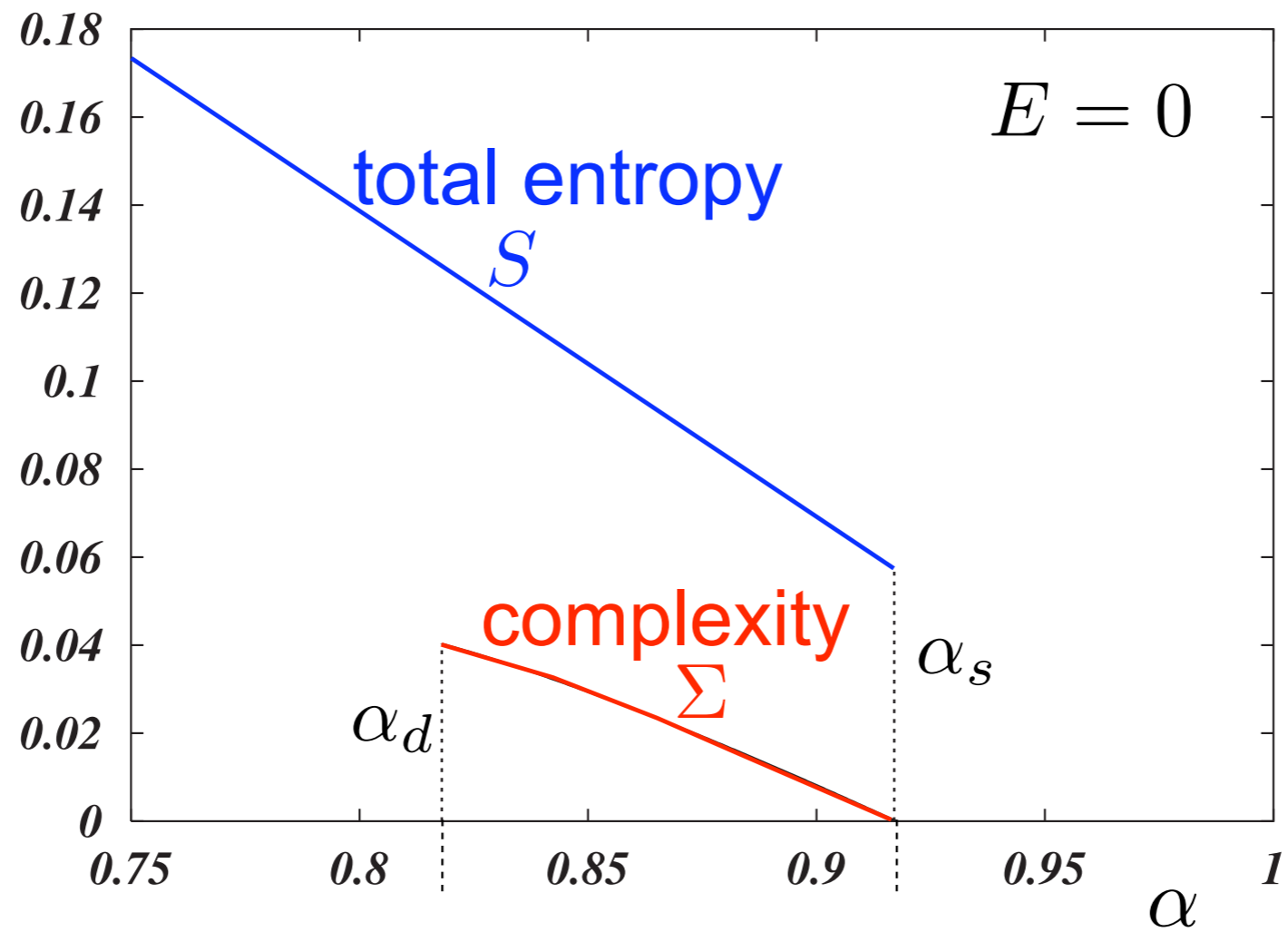
where the triplets $\langle ijk \rangle$ are randomly
chosen uniformly among $\binom{N}{3}$

It corresponds to computing the
ground state of a 3SG model

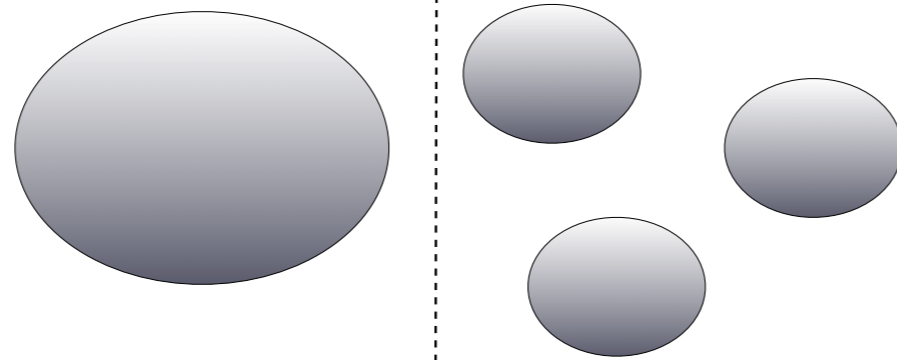
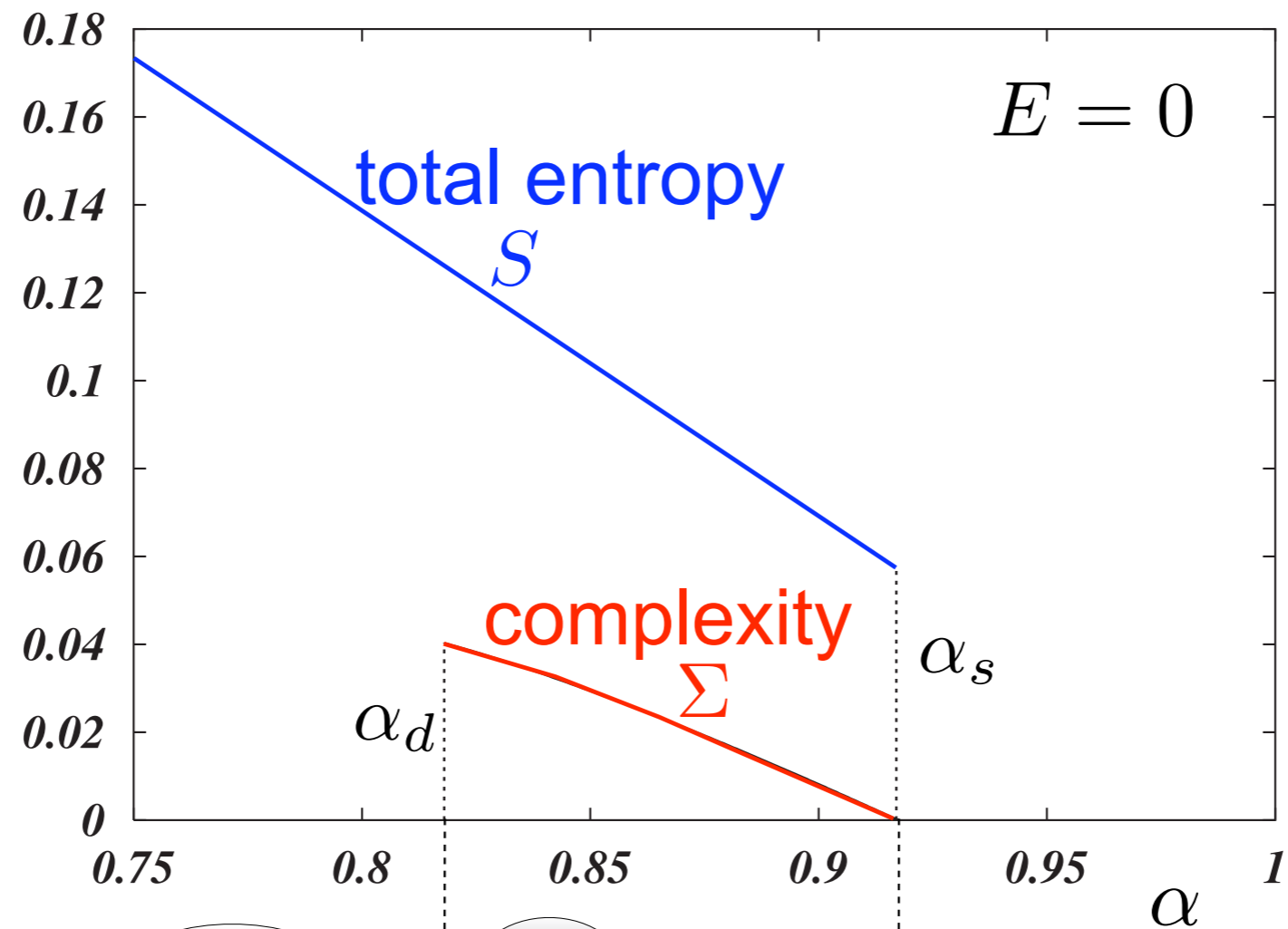
3-spin glass on a random graph



Random 3-xorsat



Random 3-xorsat

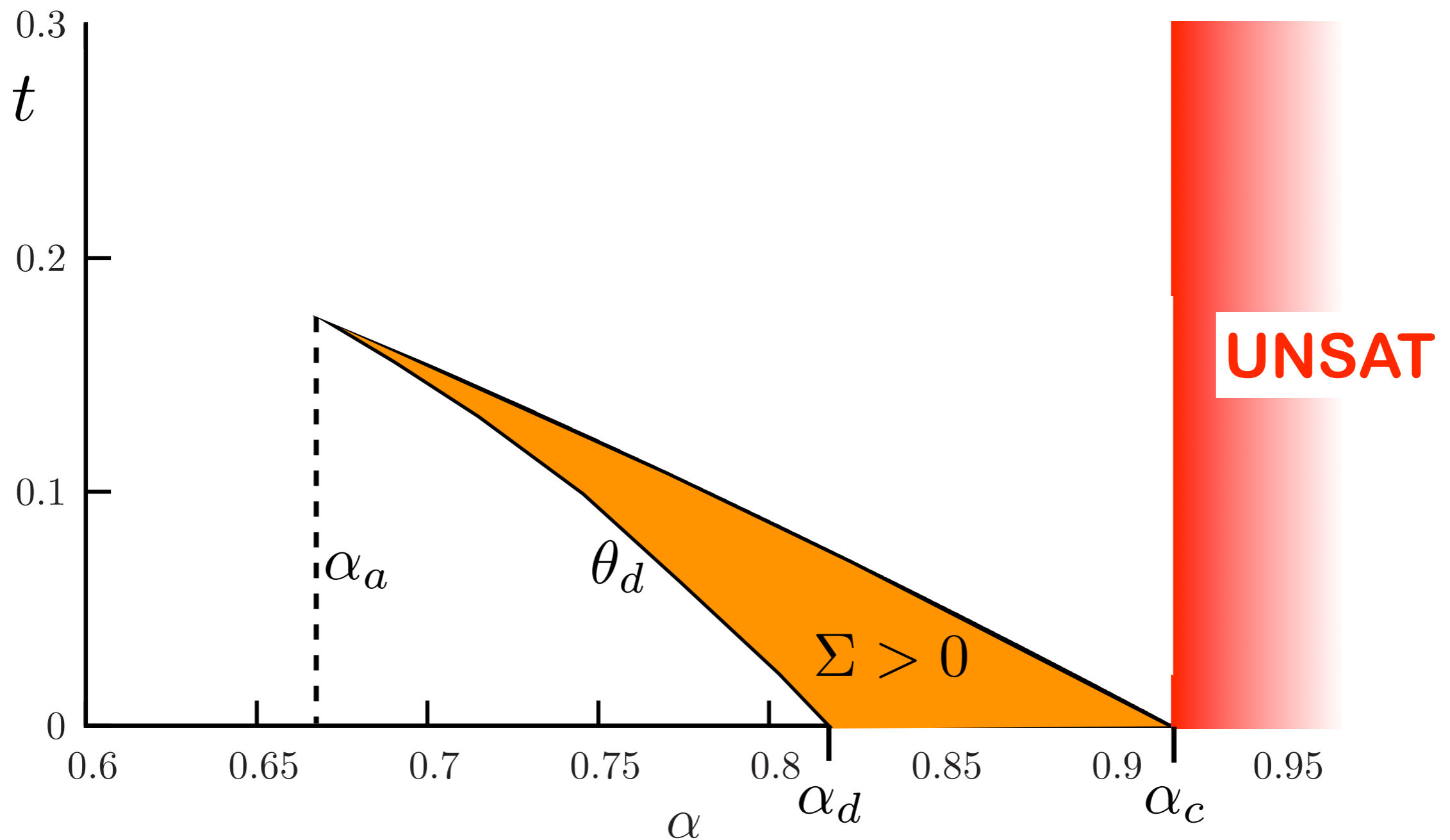


Clustering
phase transition

Dynamics: searching solutions

- Monte Carlo is very inefficient for $T \rightarrow 0$
- Sequential construction algorithm:
while (there are unassigned variables)
 - compute marginals
 - choose randomly an unassigned variable
 - fix it according to the marginal

Mapping the dynamics to a static problem...



Conclusions

- Spin glasses are prototype for complex systems
- The link between statics and dynamics has provided very useful
- ...but many aspects of this connection are still unclear and need to be improved.

Bibliography

- Too long to fit in a single slide ;-)
- Depends on which specific issue you are interested in
- Ask me via e-mail
`Federico.Ricci@roma1.infn.it`

Phase transitions in Random SAT

