# Efficient sources for spectroscopy of excited mesons in lattice QCD

#### Daniel Mohler

Inst. f. Physik, FB Theoretische Physik Universität Graz

> Wien, November 29 2008



Advisors: Christof Gattringer, Christian Lang Co-Workers: Georg Engel, Markus Limmer Leonid Glozman, Sasa Prelovsek



# Outline

#### Excited states in Lattice QCD

- Motivation and Introduction
- Variational method and interpolating fields
- Construction of efficient quark sources

#### 2 Spectroscopy with derivative sources

- Excited pions with derivative sources
- Quenched results for the (isovector) scalar mesons

#### 3 Tetraquark interpolators for scalar mesons

#### 4 Summary & Outlook

# Motivation: Mesons from quenched QCD



Burch et al., PRD 73 (2006) 094505

- Excited states are a lot more difficult than ground states!
- Goal: Improving the method for extracting excited states.

## **Euclidean correlators**

Euclidean correlator of two Hilbert-space operators Ô<sub>1</sub> and Ô<sub>2</sub>.

$$\left< \hat{O}_2(t) \hat{O}_1(0) \right> = \sum_n e^{-t E_n} < 0 |\hat{O}_2|n> < n |\hat{O}_1|0>$$

• Can also be expressed as a Euclidean path integral

$$\left\langle \hat{O}_{2}(t)\hat{O}_{1}(0)
ight
angle =rac{1}{Z}\int\mathcal{D}[\psi,ar{\psi},U]e^{-S_{E}}O_{2}[\psi,ar{\psi},U]O_{1}[\psi,ar{\psi},U]$$

- Will show how to apply this to extract the spectrum of mesons from lattice QCD.
- The same applies to other quantities of interest: baryon correlators, three point functions, matrix elements, ...

A (10) A (10) A (10)

## **Euclidean correlators**

• Euclidean correlator of two Hilbert-space operators  $\hat{O}_1$  and  $\hat{O}_2$ .

$$\left< \hat{O}_2(t) \hat{O}_1(0) \right> = \sum_n e^{-t E_n} < 0 |\hat{O}_2|n> < n |\hat{O}_1|0>$$

• Can also be expressed as a Euclidean path integral

$$\left\langle \hat{\mathsf{O}}_{2}(t)\hat{\mathsf{O}}_{1}(0)
ight
angle =rac{1}{Z}\int\mathcal{D}[\psi,ar{\psi},U]e^{-\mathcal{S}_{E}}\mathsf{O}_{2}[\psi,ar{\psi},U]\mathsf{O}_{1}[\psi,ar{\psi},U]$$

- Will show how to apply this to extract the spectrum of mesons from lattice QCD.
- The same applies to other quantities of interest: baryon correlators, three point functions, matrix elements, ...

# Quenched approximation vs. full QCD

"Full QCD":

$$C(t) \propto \int \mathcal{D}[U] \mathcal{D}[\psi, \bar{\psi}] e^{-S_G[U] - \bar{\psi} D[U] \psi} M(t) \bar{M}(0)$$
  
=  $\int \mathcal{D}[U] e^{-S_G[U]} (\det D_u \det D_d \dots)$   
 $\times \left[ D_u^{-1} D_d^{-1} \dots + \dots \right]$ 

- Set det D ≡ 1 (no dynamical fermion vacuum, i.e. no sea quarks)
- Gauge field vacuum is fully dynamical (Monte Carlo)
- Consider only the valence quarks
- Hadron correlation functions are built from the quark propagators



Daniel Mohler (KFUGPhysics)

# Quenched approximation vs. full QCD

Quenched approximation:

$$\begin{split} \mathcal{C}(t) &\propto \int \mathcal{D}[U] \ \mathcal{D}[\psi, \bar{\psi}] \ \mathrm{e}^{-\mathcal{S}_{\mathsf{G}}[U] - \bar{\psi} \ \mathcal{D}[U] \ \psi} \ \mathcal{N}(t) \bar{\mathcal{N}}(0) \\ &= \int \mathcal{D}[U] \ \mathrm{e}^{-\mathcal{S}_{\mathsf{G}}[U]} \\ &\times \left[ \mathcal{D}_{u}^{-1} \mathcal{D}_{d}^{-1} \dots + \dots \right] \end{split}$$

- Set det D ≡ 1 (no dynamical fermion vacuum, i.e. no sea quarks)
- Gauge field vacuum is fully dynamical (Monte Carlo)
- Consider only the valence quarks
- Hadron correlation functions are built from the quark propagators



Variational method (C.Michael; Lüscher and Wolff) Matrix of correlators projected to fixed momentum (will assume 0)

$$C(t)_{ij} = \sum_{n} \mathrm{e}^{-tE_{n}} \left< 0 \left| O_{i} \right| n \right> \left< n \left| O_{j}^{\dagger} \right| 0 \right>$$

Solve the generalized eigenvalue problem:

$$egin{split} \mathcal{C}(t)ec{\psi}^{(k)} &= \lambda^{(k)}(t)\mathcal{C}(t_0)ec{\psi}^{(k)} \ \lambda^{(k)}(t) \propto \mathrm{e}^{-t\mathcal{E}_k}\left(1+\mathcal{O}\left(\mathrm{e}^{-t\Delta\mathcal{E}_k}
ight)
ight) \end{split}$$

At large time separation: only a single mass in each eigenvalue. Eigenvectors can serve as a fingerprint.

4 3 5 4 3 5

• Interesting observation for  $t \leq 2t_0$ 

$$\Delta E_k \equiv E_{N+1} - E_k$$

Blossier et al. [arXiv: 0808.1017]

• Can also determine couplings:

$$C(t) = \sum_{n=1}^{\infty} v_i^{(n)} v_i^{(n)*} e^{-tE^{(n)}} \text{ with } v_i^{(n)} = \left\langle 0 | O_i | H^{(n)} \right\rangle$$
$$R(t)_i^{(k)} = \frac{|\sum_j C(t)_{ij} \psi_j^{(k)}|^2}{\sum_k \sum_l \psi_k^{(k)*} C(t)_{kl} \psi_l^{(k)}} \approx v_i^{(k)} v_i^{(k)*} e^{-tE^{(k)}}$$

Burch et al. [arXiv: 0809.1103]

# Interpolating fields for mesons

- We need interpolating field operators with quantum numbers J<sup>PC</sup>
- These are typically build from smeared quarks

 $\bar{\psi} \mathsf{\Gamma} \psi$ 

• Some examples (for Isovector mesons):

JPC	interpolator	J <sup>PC</sup>	interpolator
0-+	$\bar{u}_{ m s}\gamma_5 u_{ m s}$	0-+	$\bar{u}_{s}\partial_{i}\gamma_{i}\gamma_{t}\gamma_{5}u_{s}$
1	$ar{u}_{s}\gamma_{i}u_{s}$	1	ū <sub>s</sub> ∂ <sub>i</sub> u <sub>s</sub>

• Remark: At finite lattice spacing *a* interpolators will also couple to higher spin states.

A (10) > A (10) > A (10)

## Interpolating fields for mesons

- We need interpolating field operators with quantum numbers J<sup>PC</sup>
- These are typically build from smeared quarks

 $\bar{\psi} \Gamma \psi$ 

• Some examples (for Isovector mesons):

J <sup>PC</sup>	interpolator	· · · · ·	J <sup>PC</sup>	interpolator
0-+	$ar{u}_{s}\gamma_{5}u_{s}$		0-+	$ar{u}_{s}\partial_{i}\gamma_{i}\gamma_{t}\gamma_{5}u_{s}$
1	$ar{u}_{s}\gamma_{i}u_{s}$		1	$\bar{u}_{s}\partial_{i}u_{s}$

• Remark: At finite lattice spacing *a* interpolators will also couple to higher spin states.

# Interpolating fields for mesons

- We need interpolating field operators with quantum numbers J<sup>PC</sup>
- These are typically build from smeared quarks

 $\bar{\psi} \Gamma \psi$ 

• Some examples (for Isovector mesons):

J <sup>PC</sup>	interpolator	$J^{PC}$	interpolator
0-+	$ar{u}_{s}\gamma_{5}u_{s}$	0-+	$\bar{u}_{s}\partial_{i}\gamma_{i}\gamma_{t}\gamma_{5}u_{s}$
1	$ar{u}_{ m s}\gamma_i u_{ m s}$	1	$\bar{u}_{s}\partial_{i}u_{s}$

 Remark: At finite lattice spacing a interpolators will also couple to higher spin states.

A (10) > A (10) > A (10)

#### Quark sources

• Jacobi smeared quark sources, e.g.,  $u_s \equiv (S u)_x$ 

$$S = M S_0 \quad \text{with} \quad M = \sum_{n=0}^{N} \kappa^n H^n$$
$$H(\vec{n}, \vec{m}) = \sum_{j=1}^{3} \left( U_j(\vec{n}, 0) \delta\left(\vec{n} + \hat{j}, \vec{m}\right) + U_j\left(\vec{n} - \hat{j}, 0\right)^{\dagger} \delta\left(\vec{n} - \hat{j}, \vec{m}\right) \right)$$

Combination allows nodes in the interpolating operators



Derivative quark sources  $W_{d_i}$ :

$$D_i(\vec{x}, \vec{y}) = U_i(\vec{x}, 0)\delta(\vec{x} + \hat{i}, \vec{y}) - U_i(\vec{x} - \hat{i}, 0)^{\dagger}\delta(\vec{x} - \hat{i}, \vec{y}) ,$$
  
$$W_{d_i} = D_i S .$$



#### Mesons with derivative sources: Pion channel



Figure: 1st and 2nd excitation of  $\pi$ 

Gattringer et al., ArXiv:0802.2020 [hep-lat], PRD 78, 034501, 2008



Figure: Eigenvector components for ground state and lowest excitations

Gattringer et al., ArXiv:0802.2020 [hep-lat], PRD 78, 034501, 2008

#### Mesons with derivative sources: Isovector Scalar 10<sup>++</sup>



Figure: Diagonal correlators and ground state in the 0<sup>++</sup> channel

Gattringer et al., ArXiv:0802.2020 [hep-lat], PRD 78, 034501, 2008

Daniel Mohler (KFUGPhysics)

Excited mesons in lattice QCD

Mesons with derivative sources: Isovector Scalar 10<sup>++</sup>



Figure: Diagonal correlators and ground state in the 0<sup>++</sup> channel

Gattringer et al., ArXiv:0802.2020 [hep-lat], PRD 78, 034501, 2008

Other quenched results (also) extrapolate to  $a_0(1450)$ 

Daniel Mohler (KFUGPhysics)

Excited mesons in lattice QCD

# The scalar meson puzzle

#### Low lying scalars could be tetraquark states



- quark models would place  $\bar{q}q$  with L = 1 above 1GeV
- $m_{\kappa} < m_{a_0}$  hard to reconcile with  $\bar{u}s$  and  $\bar{u}d$
- $a_0(980)$  couples well with  $K\bar{K}$

**H N** 

Tetraquarks with a diquark - anti-diquark structure

$$\begin{split} [q\mathbf{Q}]_{a} &\equiv \epsilon_{abc} [q_{b}^{T} C \gamma_{5} \mathbf{Q}_{c} - \mathbf{Q}_{b}^{T} C \gamma_{5} q_{c}] \\ [\bar{q}\bar{\mathbf{Q}}]_{a} &\equiv \epsilon_{abc} [\bar{q}_{b} C \gamma_{5} \bar{\mathbf{Q}}_{c}^{T} - \bar{\mathbf{Q}}_{b} C \gamma_{5} \bar{q}_{c}^{T}] , \end{split}$$

We simulate

$$\mathcal{O}^{l=0} = [\textit{ud}][\bar{\textit{u}}\bar{\textit{d}}] \;, \quad \mathcal{O}^{l=1/2} = [\textit{ud}][\bar{\textit{d}}\bar{\textit{s}}] \;, \quad \mathcal{O}^{l=1} = [\textit{us}][\bar{\textit{d}}\bar{\textit{s}}]$$

with different source/sink smearings and use the variational method. S.Prelovsek and D.M., arXiv:0810.1759

#### Scalar tetraquarks - results



Main conclusion: No evidence of low lying scalar tetraquarks for  $M_{\pi}$  in the range 344...576 MeV

S.Prelovsek and D.M., arXiv:0810.1759

- Excited states are difficult on the lattice.
- Variational method can serve to reliably extract excited states.
- Presented results from an improved basis containing interpolators with derivatives.
- Presented results for scalar mesons using tetraquark interpolators.
- Analysis of dynamical data is ongoing.
- Improved basis can also be used for baryons and tetraquarks.
- Scattering states and resonances further complicate the analysis.
- Want to apply methods to charmonium states using MILC lattices (With C. DeTar and T. Burch).

- Excited states are difficult on the lattice.
- Variational method can serve to reliably extract excited states.
- Presented results from an improved basis containing interpolators with derivatives.
- Presented results for scalar mesons using tetraquark interpolators.
- Analysis of dynamical data is ongoing.
- Improved basis can also be used for baryons and tetraquarks.
- Scattering states and resonances further complicate the analysis.
- Want to apply methods to charmonium states using MILC lattices (With C. DeTar and T. Burch).

A (10) A (10) A (10)