

# Symbolic Derivation of Dyson-Schwinger Equations using *Mathematica*

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# The Quark Propagator Dyson-Schwinger Equation

DSEs describe the non-perturbative propagation  
and interaction of particles.

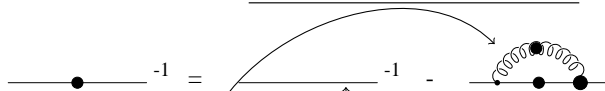
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$$\text{---} \bullet \text{---}^{-1} = \text{---} \text{---}^{-1} - \text{---} \bullet \text{---} \text{---} \bullet \text{---} \bullet \text{---}$$

# The Quark Propagator Dyson-Schwinger Equation

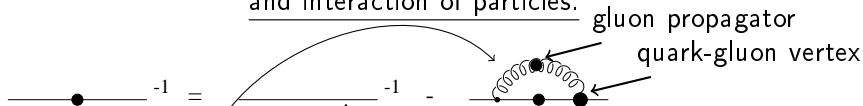
DSEs describe the non-perturbative propagation and interaction of particles.



- Propagation of a quark: nothing happens or it emits/absorbs a gluon "in all possible ways" → dressed propagators and vertices

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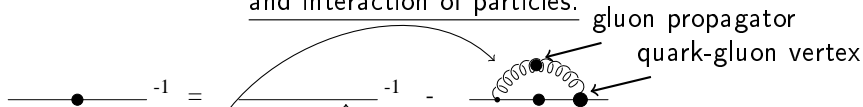
DSEs describe the non-perturbative propagation and interaction of particles.



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- **Unknowns:** Quark-gluon vertex, gluon propagator

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- **Unknowns: Quark-gluon vertex, gluon propagator**
- Spontaneous chiral symmetry breaking  
⇒ dynamically generated momentum dependent quark masses

# Dyson-Schwinger Equations (DSEs) for Investigating QCD

## Facts about DSEs

- F. J. Dyson (1949) and J. S. Schwinger (1951)
- Equations of motion of Green functions
- Infinitely large tower of equations (DSE for  $n$ -point function contains  $n+1$ - and  $n+2$ -point functions)



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→ non-perturbative regime accessible
- Continuum  
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Pros:

- Exact equations  
→ **non-perturbative regime** accessible
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Cons?:

- Truncations (not for all tasks)
- **Gauge-dependent**  
→ Exploit advantages of different gauges

## How to Derive a DSE? (I)

$\phi$ : **generic field** (quarks, gluons, ghosts, scalars, ...).

Put all indices (Lorentz, color), coordinates and field type in one index.  
Integration and summation understood.

Start with the **path integral**. Integral of a total derivative vanishes (given appropriate boundary conditions):

$$Z[J] = \int D[\phi] e^{-\mathbf{S}[\phi] + J\phi} \longrightarrow \frac{\delta Z}{\delta \phi_i} = \left( - \frac{\delta \mathbf{S}}{\delta \phi'_i} \Big|_{\phi'_i = \delta / \delta J_i} + J_i \right) Z[J] = 0$$

$$- \frac{\delta \mathbf{S}}{\delta \phi'_i} \Big|_{\phi'_i = \delta W / \delta J_i + \delta / \delta J_i} + J_i = 0 \quad \text{connected n-point functions}$$

$$- \frac{\delta \mathbf{S}}{\delta \phi'_i} \Big|_{\phi'_i = \phi_i + D_{ij} \delta / \delta \phi_j} + \frac{\delta \Gamma}{\delta \phi_i} = 0 \quad \text{1PI n-point functions}$$

## How to Derive a DSE? (II)

Continue by **applying derivatives** with respect to the fields:

$$\frac{\delta^n \Gamma}{\delta \phi_{i_1 \dots i_n}} = \frac{\delta \mathcal{S}}{\delta^n \phi'_{i_1 \dots i_n}} \Big|_{\phi'_i = \phi_i + D_{ij} \delta / \delta \phi_j}$$

$\phi = 0$  only at the end

$\Rightarrow$  field-dependent **intermediate Green functions**.

Number of terms grows with  $n$  and number of interactions!

1 field, 3- and 4-point interactions:

	# diagrams
2-point	5
3-point	12
4-point	53
5-point	359
...	...

Automatization possible?



# Diagrammatic Rules

What does applying a derivative mean

for a field?  $\frac{\delta\phi_j}{\delta\phi_i} = \delta_{ij}$

$$\frac{\delta}{\delta\phi_i} \parallel = \parallel_i \otimes$$

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for a vertex?  $\frac{\delta}{\delta \phi_i} \Gamma_{j_1 \dots j_n} = \frac{-\delta^{n+1} \Gamma}{\delta \phi_i \delta \phi_{j_1} \dots \delta \phi_{j_n}} = \Gamma_{ij_1 \dots j_n}$

$$\frac{\delta}{\delta \phi_i} \bullet = \bullet_i$$



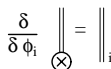
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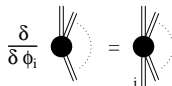
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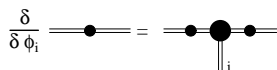
for a field?  $\frac{\delta \phi_j}{\delta \phi_i} = \delta_{ij}$

for a vertex?  $\frac{\delta}{\delta \phi_i} \Gamma_{j_1 \dots j_n} = \frac{-\delta^{n+1} \Gamma}{\delta \phi_i \delta \phi_{j_1} \dots \delta \phi_{j_n}} = \Gamma_{ij_1 \dots j_n}$

for a propagator?  $\frac{\delta}{\delta \phi_i} D_{jk} = \frac{\delta}{\delta \phi_i} \left( \frac{\delta^2 \Gamma}{\delta \phi_j \delta \phi_k} \right)^{-1} = D_{jj'} \Gamma_{j'ik'} D_{k'k}$

$$\frac{\delta}{\delta \phi_i} \parallel = \parallel_i$$


$$\frac{\delta}{\delta \phi_i} \bullet = \bullet_i$$


$$\frac{\delta}{\delta \phi_i} \text{---} \bullet = \text{---} \bullet \text{---} \bullet_i$$


# Diagrammatic Rules

What does applying a derivative mean

for a field?

$$\frac{\delta \phi_j}{\delta \phi_i} = \delta_{ij}$$

$$\frac{\delta}{\delta \phi_i} \text{ (line with } \otimes \text{)} = \text{ (line with } \parallel \text{)}_i$$

for a vertex?

$$\frac{\delta}{\delta \phi_i} \Gamma_{j_1 \dots j_n} = \frac{-\delta^{n+1} \Gamma}{\delta \phi_i \delta \phi_{j_1} \dots \delta \phi_{j_n}} = \Gamma_{ij_1 \dots j_n}$$

$$\frac{\delta}{\delta \phi_i} \text{ (n lines meeting at a dot)} = \text{ (n lines meeting at a dot with } \parallel \text{)}_i$$

for a propagator?

$$\frac{\delta}{\delta \phi_i} D_{jk} = \frac{\delta}{\delta \phi_i} \left( \frac{\delta^2 \Gamma}{\delta \phi_j \delta \phi_k} \right)^{-1} = D_{ij'} \Gamma_{j'ik'} D_{k'k}$$

$$\frac{\delta}{\delta \phi_i} \text{ (double line)} = \text{ (double line with } \bullet \text{)}_i$$

Ex.: Derivative with respect to gluon in Landau gauge QCD

$$\frac{\delta}{\delta A_i} \text{ (gluon line with } \otimes \text{)} = \text{ (gluon line with } \parallel \text{)}_i$$

$$\frac{\delta}{\delta A_i} \text{ (gluon line with } \bullet \text{)} = \text{ (gluon line with } \bullet \text{ and } \parallel \text{)}_i$$

$$\frac{\delta}{\delta A_i} \text{ (gluon loop)} = \text{ (gluon loop with } \parallel \text{)}_i$$

$$\frac{\delta}{\delta A_i} \text{ (gluon vertex)} = \text{ (gluon vertex with } \parallel \text{)}_i$$

$$\frac{\delta}{\delta A_i} \text{ (gluon propagator)} = \text{ (gluon propagator with } \bullet \text{)}_i$$

generic field!

## Example: Landau Gauge Yang-Mills theory

$$\mathcal{L} = \frac{1}{4} F_{\mu\nu}^a F_{\mu\nu}^a$$

$$F_{\mu\nu}^a = \partial_\mu A_\nu^a - \partial_\nu A_\mu^a - g f^{abc} A_\mu^b A_\nu^c$$

Gluonic self-interaction!

Gauge fixed to Landau gauge.

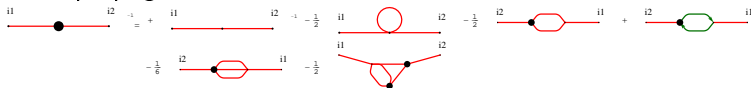
- 2 propagators: gluons  $A$  and ghosts  $c$
- 3 interactions:  $AAA$ ,  $Acc$ ;  $AAAA$

We start from

$$\frac{\delta^n \Gamma}{\delta \phi_{i_1 \dots i_n}} = \frac{\delta \mathcal{S}}{\delta \phi'_{i_1 \dots i_n}} \Big|_{\phi'_i = \phi_i + D_{ij} \delta / \delta \phi_j}$$

# Landau Gauge: Propagators

**Gluon** propagator:



**Ghost** propagator:



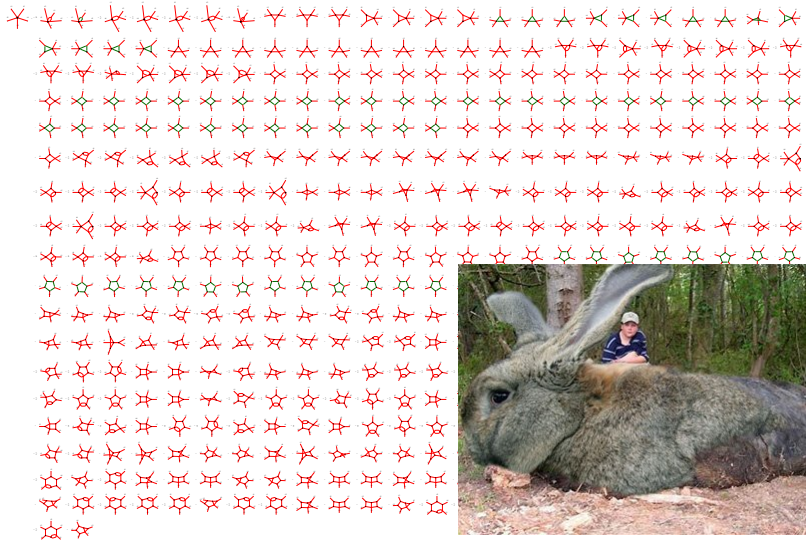
# Landau Gauge: Four-Gluon Vertex

66 terms



# Landau Gauge: Five-Gluon Vertex

434 terms



## More Fields and Interactions

Now imagine you have 3 fields and 11 interactions.



### Yang-Mills theory in the maximally Abelian gauge

- 3 fields: diagonal gluons  $A$ , off-diagonal gluons  $B$ , ghosts  $c$
- 4 three-point interactions:  $ABB$ ,  $A_{cc}$ ;  $BBB$ ,  $B_{cc}$
- 7 four-point interactions:  $AABB$ ,  $AA_{cc}$ ,  $BBBB$ ,  $BB_{cc}$ ,  $cccc$ ;  $ABBB$ ,  $AB_{cc}$

If you have only one field, every type of graph appears once.

If you have several fields, **every type of graph** can appear in **several variations**.

→ **Symmetries of the Lagrangian**, e. g. ghost-number conservation, restriction in number of diagonal gluons.



# The Infrared Solution of the Maximally Abelian Gauge

## Non-linear gauge fixing

Minimize off-diagonal components:

$$D_{\mu}^{ab} B^b = (\delta_{ab} \partial_{\mu} - g f^{abi} A_{\mu}^i) B^b = 0$$

Scaling solution: **power laws** for dressing functions in the IR, e. g.

$$\Delta_{\mu\nu}^A = (p^2)^{\delta_A - 1} P_{\mu\nu},$$

with **exponents**

$$-\delta_A = \delta_B = \delta_c \geq 0.$$

[R. A., M. Q. H., K. S., in prep.]

⇒ Diagonal gluon relevant degree of freedom = infrared enhanced.

**Unique solution** for all vertices with even number of legs.

No truncations!

(Landau gauge: ghost  $\cong$  relevant d.o.f)





# DoDSE: Derivation of Dyson-Schwinger Equations

- Idea: **Automatize the derivation process** of DSEs.
- Use algorithm explained above in symbolic programming language  $\rightarrow$  *Mathematica*.
- Draw corresponding **Feynman diagrams**?

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## Structures in *DoDSE*

- Fields: list of name and index
- Vertices:  $V$
- Propagators:  $P$
- Bare quantities:  $S$
- Diagrams:  $op$

This is the minimum information needed to write down one expression unambiguously. Still...



# Output of *DoDSE*

$$\text{op}[S[\{A, i1\}, \{A, i2\}]]$$

$$-(1/2) \text{op}[S[\{A, i1\}, \{A, i2\}, \{A, r1\}, \{A, s1\}], P[\{A, r1\}, \{A, s1\}]]$$

$$-1/2 \text{op}[S[\{A, i1\}, \{A, r1\}, \{A, s1\}], V[\{A, i2\}, \{A, t1\}, \{A, u1\}], P[\{A, r1\}, \{A, t1\}], P[\{A, s1\}, \{A, u1\}]]$$

$$\text{op}[S[\{A, i1\}, \{cb, r1\}, \{c, s1\}], V[\{A, i2\}, \{cb, u1\}, \{c, t1\}], P[\{c, t1\}, \{cb, r1\}], P[\{c, s1\}, \{cb, u1\}]]$$

$$-(1/6) \text{op}[S[\{A, i1\}, \{A, r1\}, \{A, r2\}, \{A, s1\}], P[\{A, r1\}, \{A, s2\}], P[\{A, r2\}, \{A, t2\}], P[\{A, s1\}, \{A, u2\}], V[\{A, i2\}, \{A, s2\}, \{A, t2\}, \{A, u2\}]]$$

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Fields, indices, propagators, vertices

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# Output of *DoDSE*

$$\begin{aligned}
 &+ \text{il} \text{---} \text{i2} \quad -1 \\
 &\text{op}[S[\{A, \text{il}\}, \{A, \text{i2}\}]] \\
 \\
 &-\frac{1}{2} \text{il} \text{---} \text{i2} \\
 &-(1/2) \text{op}[S[\{A, \text{il}\}, \{A, \text{i2}\}, \{A, \text{r1}\}, \{A, \text{s1}\}], P[\{A, \text{r1}\}, \{A, \text{s1}\}]] \\
 \\
 &-\frac{1}{2} \text{i2} \text{---} \text{il} \\
 &-1/2 \text{op}[S[\{A, \text{il}\}, \{A, \text{r1}\}, \{A, \text{s1}\}], V[\{A, \text{i2}\}, \{A, \text{t1}\}, \{A, \text{u1}\}], \\
 &\quad P[\{A, \text{r1}\}, \{A, \text{t1}\}], P[\{A, \text{s1}\}, \{A, \text{u1}\}]] \\
 \\
 &+ \text{i2} \text{---} \text{il} \\
 &\text{op}[S[\{A, \text{il}\}, \{\text{cb}, \text{r1}\}, \{c, \text{s1}\}], V[\{A, \text{i2}\}, \{\text{cb}, \text{u1}\}, \{c, \text{t1}\}], \\
 &\quad P[\{c, \text{t1}\}, \{\text{cb}, \text{r1}\}], P[\{c, \text{s1}\}, \{\text{cb}, \text{u1}\}]] \\
 \\
 &-\frac{1}{6} \text{i2} \text{---} \text{il} \\
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 \end{aligned}$$

Fields, indices, propagators, vertices

## Basic Input for *DoDSE*

Fields, propagators, interactions, e. g.

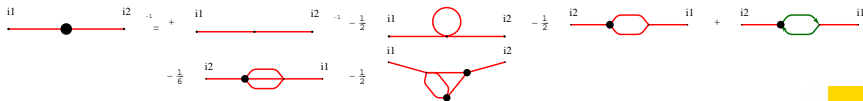
$IA = \{\{A, A\}, \{cb, c\}, \{A, cb, c\}, \{A, A, A\}, \{A, A, A, A\}\}$

No explicit Lagrangian!

$\Rightarrow AA = \text{doDSE}[IA, \{A, A\}]$

For drawing Feynman graphs: [styles for the fields](#)

$\Rightarrow \text{DSEPlot}[AA, IA, \{\{A, \text{Red}\}, \{c, \text{Green}\}\}]$



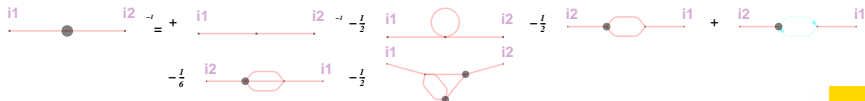
## Optional Arguments in *DoDSE*

Some **symmetries** are taken into account automatically, e. g. Grassmann number conservation.

Sometimes additional information necessary:

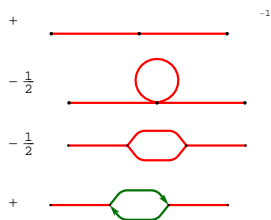
- Restrictions on vertices (e.g. maximally Abelian gauge)  $\rightarrow$  done with simple function.
- Mixed propagators  $\rightarrow$  provide list of allowed possibilities.

For drawing Feynman graphs: **options** for indices, numerical factors and output format (equation/list)



# Outlook: *Symb2Alg*


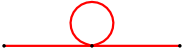


Symbolic expressions are fine for some tasks, e. g. infrared analysis, but also **algebraic expressions** are needed!





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+		-1	$\frac{(g^{\mu_1 \nu_1} p_1^2 - p_1^{\mu_1} p_1^{\nu_1}) \delta_{a_1 b_1}}{p_1^2}$
- $\frac{1}{2}$			$-\frac{C_A g^2 (q_1^{\mu_1} q_1^{\nu_1} + 2 g^{\mu_1 \nu_1} q_1^2) \delta_{a_1 b_1}}{q_1^2}$
- $\frac{1}{2}$			$-\frac{C_A g^2 (\ll 1 \gg) \delta_{a_1 b_1}}{2 \ll 2 \gg^2 (p_1^2 + 2 p_1 \cdot \ll 2 \gg + q_1^2)^2}$
+			$\frac{C_A g^2 q_1^{\mu_1} (p_1^{\nu_1} + q_1^{\nu_1}) \delta_{a_1 b_1}}{q_1^2 (p_1^2 + 2 p_1 \cdot q_1 + q_1^2)}$

*Mathematica* package *Symb2Alg*: Transforms output of *DoDSE* into algebraic expressions.

Depending on Feynman rules compatible with *FeynCalc*.

# Summary

- DSEs are important tools in some **quantum field theories**.
- Derivation of DSEs is straightforward but tedious.
- Simplification in form of **diagrammatic rules**.
- *DoDSE* implements this algorithm into a *Mathematica* package.

## *DoDSE* provides

- a convenient way to **derive DSEs**.
- basic possibilities to **draw DSEs**.
- output that can directly be used in further calculations, e. g. for infrared analysis or to get algebraic expressions that can be manipulated with *FeynCalc* or other functions.