

Light Hadron Masses from First Principles

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5th Vienna Central European Seminar
on Particle Physics and Quantum Field Theory

Outline

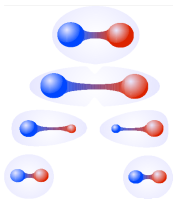
- 1 Introduction
 - Reminder: QCD
 - Reminder: LQCD
 - Why the spectrum?
- 2 Simulation details
 - Fundamentals
 - Action and simulation algorithm
 - Calculating masses
- 3 Data analysis
 - Go to physical quark masses
 - Go to infinite volume
 - Go to the continuum
 - Combined data analysis
- 4 Final result

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Important properties of QCD

QCD has remarkable properties:

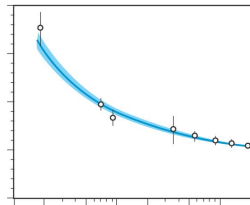


- Confinement (conformal anomaly)
- Dynamical chiral symmetry breaking of $SU_A(N_f)$ ($\langle \bar{\psi}_i \psi_i \rangle \neq 0$)
- Chiral anomaly breaking of $U_A(1)$ ($\delta \mathcal{D} \bar{\psi}_i \mathcal{D} \psi \neq 0$)

- Asymptotic freedom

$$\beta(\alpha_s) = \mu \frac{\partial \alpha_s(\mu)}{\partial \mu} = \frac{\alpha_s^2}{\pi} b_1 + \frac{\alpha_s^3}{\pi^2} b_2 + \dots$$

$$b_1 = - \left[\frac{11}{6} C_A - \frac{2}{3} \sum_R n_R T_R \right]$$



Definition of QCD

The Lagrangian of Quantum Chromodynamics (QCD)

$$\mathcal{L}_E^{QCD} = \frac{1}{2g^2} \text{Tr} [G_{\mu\nu} G_{\mu\nu}] - \sum_{i=1}^{N_f} \bar{\psi}_i (\not{D} - m_i) \psi_i + i\theta \frac{1}{16\pi^2 g^2} \text{Tr} [G_{\mu\nu} G_{\rho\sigma} \epsilon_{\mu\nu\rho\sigma}].$$

contains all possible (renormalizable) operators that are compatible with the defining symmetries of QCD:

- Poincaré symmetry
- SU(3) (color) gauge invariance
- global $SU_V(N_f) \times U_V(1) \times SU_A(N_f) \{ \times U_A(1) \}$ (massless case)
- if $\theta = 0$ C,P and T

Definition of QCD

Given the action, observables can be extracted from path-integral (PI) correlation functions

$$\begin{aligned}\langle \mathcal{O} \rangle_{QCD}^{F,G} &= \int \mathcal{D}\bar{\psi} \mathcal{D}\psi \mathcal{D}A \mathcal{O}(\bar{\psi}, \psi, A) \exp \left\{ iS^{QCD}(\bar{\psi}, \psi, A) \right\} \\ &= \int \mathcal{D}A \left(\prod_{i=1}^{N_f} \det[D(m_i)] \right) \langle \mathcal{O}(A) \rangle^F \exp \left\{ iS_G^{QCD}(A) \right\}\end{aligned}$$

- The integration of the Grassmann-variables yields the fermionic determinant.
- After performing a Wick-rotation to imaginary time, the action transforms

$$iS^{QCD} \rightarrow -S_E^{QCD}$$

and becomes *real* (\Rightarrow **LQCD simulations**)

The Lattice and the classical Wilson gauge action

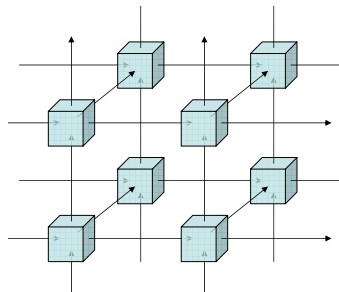
$$S_{WG} = g_0^{-2} \sum_{x, \mu\nu} \left[N_c - \Re \text{Tr} [U_\mu(x) U_\nu(x + a \hat{\mu}) U_\mu^\dagger(x + a \hat{\nu}) U_\nu^\dagger(x)] \right]$$

$$U_\mu(x) = \exp\{ia\tau^a A_\mu^a(x)\}$$

$$U_{\mu\nu}(x) = \mathbf{1} + i(a^2 G_{\mu\nu} + \mathcal{O}(a^3))$$

$$- \frac{1}{2} a^4 G_{\mu\nu}^2 + \mathcal{O}(a^5)$$

$$S_{WG} = \frac{1}{2g_0^2} \sum_x \text{Tr} [G_{\mu\nu} G_{\mu\nu}] + \mathcal{O}(a^2)$$



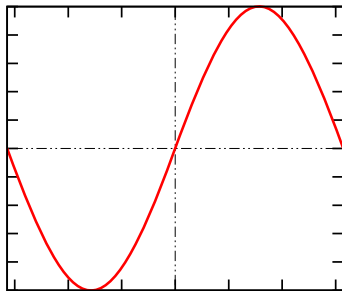
- Quarks live on the lattice points
- Gluons not only live on the links, they *are* the links

Symmetries

Discretization should preserve as many symmetries as possible.

It cannot preserve

- $SO(4)$ (\rightarrow Poincaré) invariance, but the lattice variant: the cubic group $W(4)$
- **Nielsen Ninomiya theorem:** Fermion action cannot preserve translational invariance and chiral symmetry while being real, local and bilinear



Latter problem is usually solved by **explicitly breaking** chiral symmetry in some fashion.

The Wilson fermion action

- Wilson solved doubling problem by adding a *continuum irrelevant* operator (the terms proportional to r)

$$S_{WF} = a^4 \sum_{xy} \bar{\psi}(x) M_{Wxy} \psi(y)$$

$$M_{Wxy} = \delta_{x,y} - \kappa \sum_{\mu} \left[(r - \gamma_{\mu}) U_{\mu}(x) \delta_{x,y+a\hat{\mu}} + (r + \gamma_{\mu}) U_{\mu}^{\dagger}(x) \delta_{x,y-a\hat{\mu}} \right]$$

$$M_W \propto am_0 - a\cancel{D} + \mathcal{O}(a^3) + r a^2 D^2 + \mathcal{O}(a^3)$$

- In other words by **explicitly breaking** chiral symmetry :

$$(\gamma_5 M_W + M_W \gamma_5)_{xy} = \delta_{x,y} - r \kappa \sum_{\mu} \left[U_{\mu}(x) \delta_{x,y+a\hat{\mu}} + U_{\mu}^{\dagger}(x) \delta_{x,y-a\hat{\mu}} \right]$$

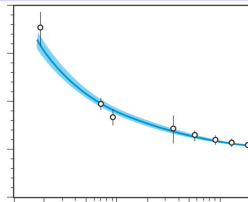
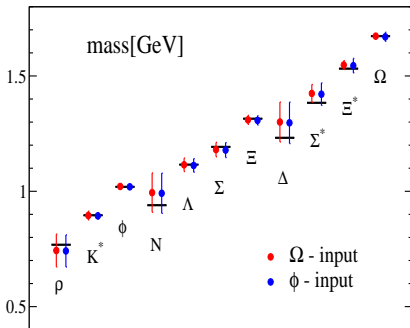
$$S_W(p) = \left[m_0 a + 4r - \sum_{\mu} (r \cos ap_{\mu} - i \gamma_{\mu} \sin ap_{\mu}) \right]$$

- Problem:** Additive quark mass renormalization, *critical slowing down*
- Improved actions & algorithms required (25 year effort by the field)

Why the spectrum?

Status

- Asymptotic freedom: good agreement between theory and experiment (perturbative methods)



- Good evidence in the non-perturbative domain (e.g. CP-PACS '07, $N_f=2+1$, $210\text{MeV} \leq M_\pi \leq 730\text{MeV}$, $a \simeq 0.087\text{ fm}$, $L \lesssim 2.8\text{ fm}$, $M_\pi L \simeq 2.9$)
- However, systematic errors **not** under control

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Importance Sampling

- After Wick-rotation, the exponent of the Boltzmann-factor in the path-integral is positive

$$\exp \left\{ iS^{QCD} \right\} \rightarrow \exp \left\{ -S_E^{QCD} \right\}$$

- The path-integral is now equivalent to a partition sum of statistical mechanics
- The Boltzmann-weight can now be interpreted as probability
- The expectation value of an operator can thus be written

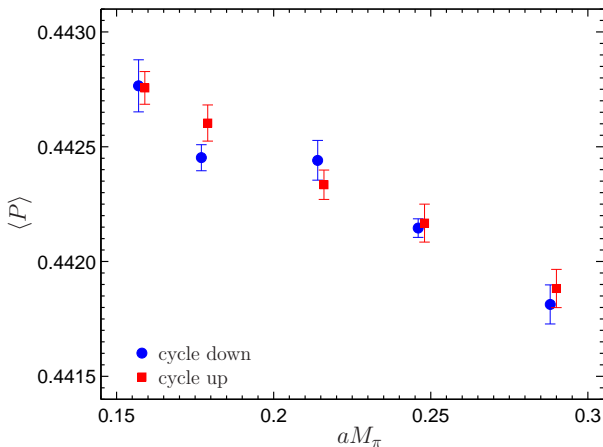
$$\langle \mathcal{O} \rangle_{F,G} = \sum_{i=0}^{\infty} \langle \mathcal{O}(A_i) \rangle_F,$$

if A_i has the probability $p \propto \exp \{ -S_E^{QCD}(A_i) \}$ to appear in the ensemble of A 's

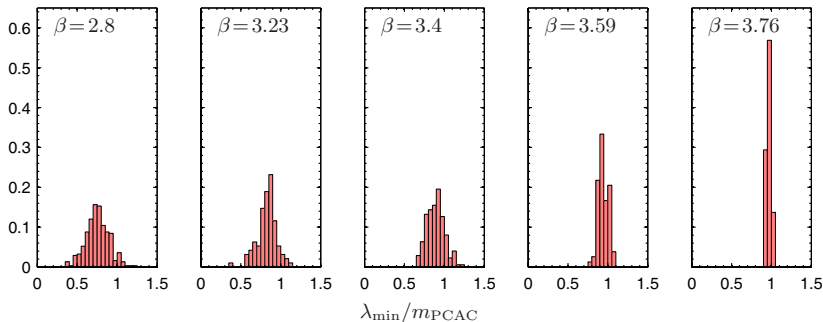
Simulation algorithm: Details

- Action
 - Clover tree level improved Wilson fermions
 - Symanzik improved gauge action
 - Stout links
- Algorithm
 - Rational HMC for strange quark
 - Mass preconditioning (“Hasenbusch trick”)
 - Multi scale integration scheme (“Sexton-Weingarten”)
 - Omelyan integrator (“non equidistant leap frog”), increasing integration precision
 - Mixed precision inverters

Algorithm stability: absence of phase transitions

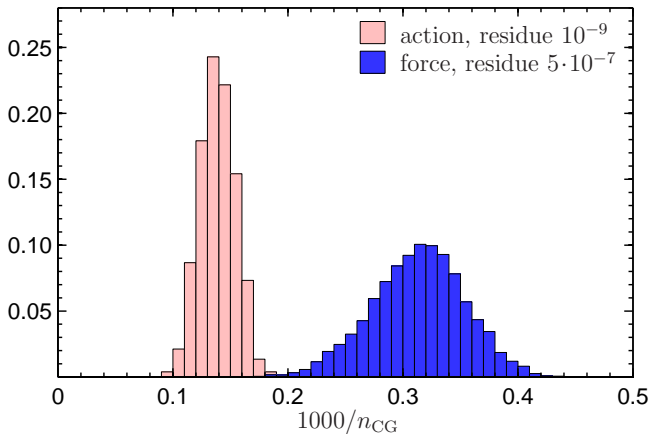


Lattices: $M_\pi \approx 240 - 440$ MeV, $a \approx 0.136$ fm

λ_{\min}^{-1} distribution

→ Simulations

Algorithm stability: distribution of $1/n_{CG}$



Lattice: $m_\pi = 190$ MeV, $a = 0.088$ fm, 64×48^3 ($m_\pi L \geq 4$)

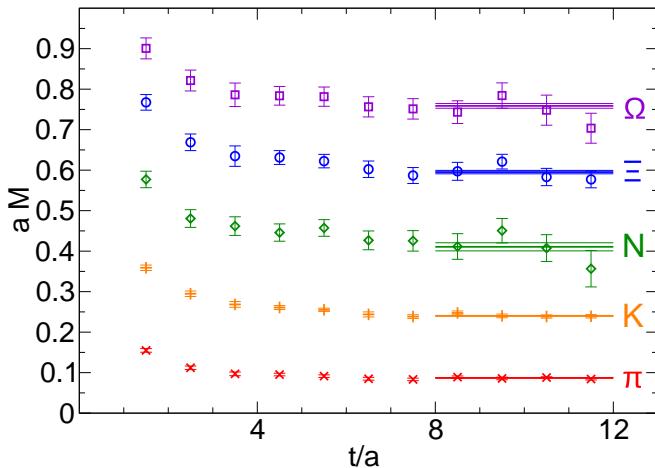
Extracting particle masses

- Masses are extracted from the formula:

$$\begin{aligned}
 \langle P(t)|P(0)\rangle &= \langle P(0)| \exp\{-\mathcal{H}t\} |P(0)\rangle \\
 &= \sum_i \frac{\langle P(0)|i\rangle \langle i|P(0)\rangle}{2E_i} \exp\{-E_i t\} \\
 &\rightarrow \frac{\langle P(0)|0\rangle \langle 0|P(0)\rangle}{2E_0} \exp\{-E_0 t\} + \mathcal{O}(\exp\{-E_1 t\}) \\
 &\stackrel{\text{BC}}{\rightarrow} \text{const.} \times \frac{\cosh}{\sinh} \{E_0 t\}
 \end{aligned}$$

- Here P will be the zero momentum projected operator

Particle correlators



Lattice: $m_\pi = 190$ MeV, $a = 0.088$ fm, 64×48^3 ($m_\pi L \geq 4$)

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Controlling **all** systematic errors

- **Make sure it's QCD:** Include u , d and s quarks into the simulation with an action, whose universality class is QCD.
- **Go to physical quark masses:** Use controlled interpolations and extrapolations of the results to physical m_{ud} and m_s
- **Go to infinite volume:** Use large volumes ($M_\pi L \gtrsim 4$) to guarantee small finite-size effects and at least one simulation at a significantly larger volume to confirm the smallness of these effects.
 - **And** treat the resonant states correctly.
- **Go to the continuum:** Use controlled extrapolations to the continuum limit, requiring that the calculations be performed at no less than three values of the lattice spacing.

Go to physical quark masses

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[Go to physical quark masses](#)

Quark mass dependence

Goal:

- Extra-/Interpolate M_X (baryon/vector meson mass) to physical point (M_π , M_K)

Method:

- Use M_Ξ or M_Ω to set the scale
- Variables to parametrize M_π^2 and M_K^2 dependence of M_X :
 - Use bare masses aM_y , $y \in \{X, \pi, K\}$ and a (bootstrapped)
 - Use dimensionless ratios $r_y := \frac{M_y}{M_{\Xi/\Omega}}$ (cancellations)

We use both procedures \rightarrow systematic error

Quark mass dependence (ctd.)

Method (ctd.):

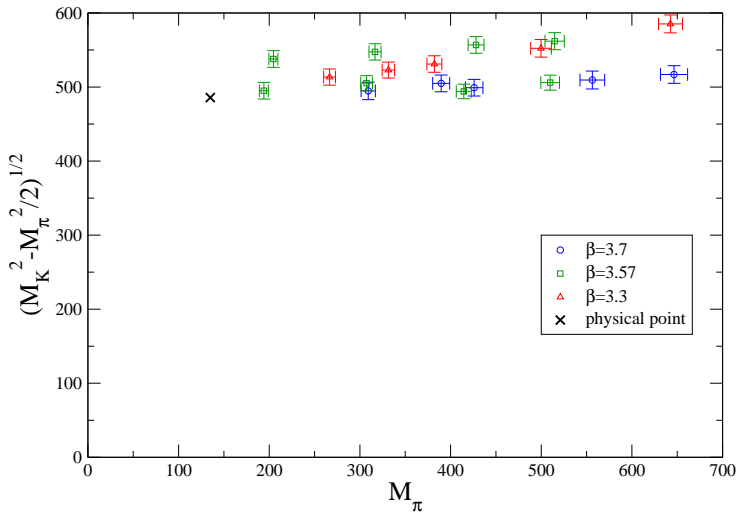
- Parametrization: $M_X = M_X^{(0)} + \alpha M_\pi^2 + \beta M_K^2 + \text{higher orders}$
 - Leading order sufficient for M_K^2 dependence
 - We include higher order term in M_π^2
 - Next order χ PT (around $M_\pi^2 = 0$): $\propto M_\pi^3$
 - Taylor expansion (around $M_\pi^2 \neq 0$): $\propto M_\pi^4$

Both procedures fine \rightarrow systematic error
No sensitivity to any order beyond these
- Vector mesons: higher orders not significant
- Baryons: higher orders significant
 - Restrict fit range to further estimate systematics:
 - Use full range, $M_\pi < 550, 450\text{MeV}$

We use all 3 ranges \rightarrow systematic error

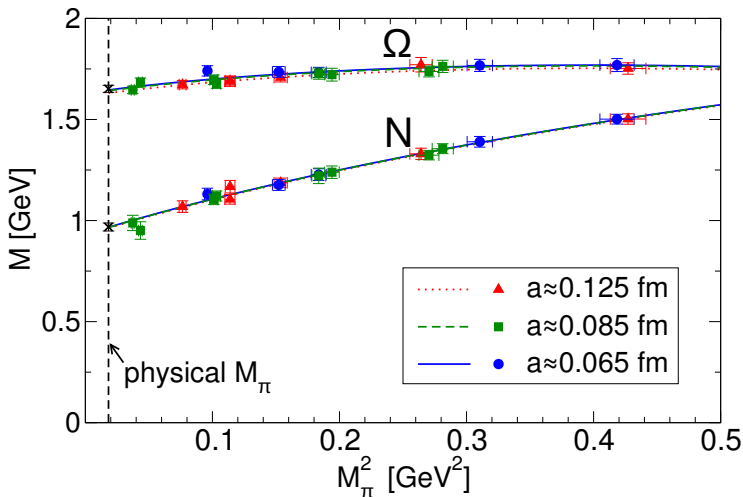
[Go to physical quark masses](#)

Data set



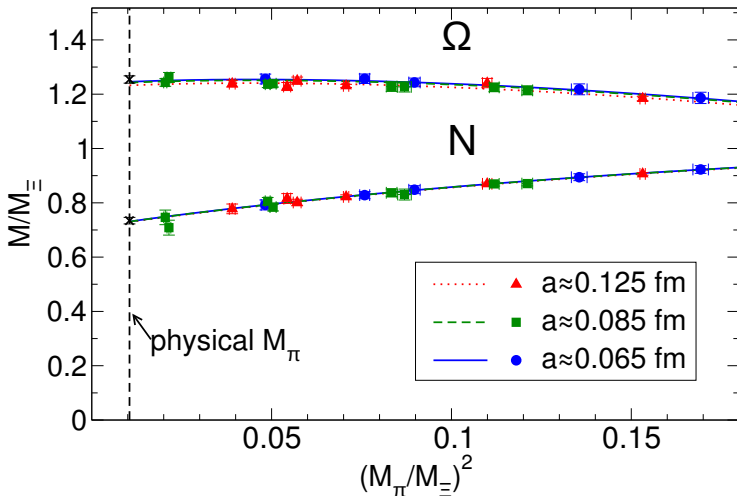
Go to physical quark masses

Chiral fit



Go to physical quark masses

Chiral fit using ratios



Go to infinite volume

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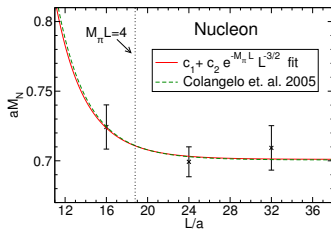
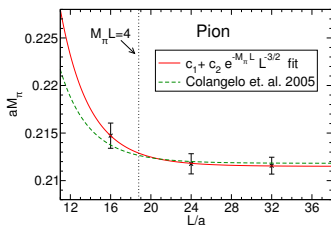
Finite volume effects from virtual pions

Goal:

- Eliminate virtual pion finite V effects

Method:

- Best practice: use large V
 - We use $M_\pi L \gtrsim 4$ (and one point to study finite V)
 - Effects are tiny and well described by $\frac{M_X(L) - M_X}{M_X} = cM_\pi^{1/2} L^{-3/2} e^{M_\pi L}$
Colangelo et. al., 2005



Finite volume effects in resonances

Goal:

- Eliminate spectrum distortions from resonances mixing with scattering states

Method:

- Stay in region where resonance is ground state
 - Otherwise no sensitivity to resonance mass in ground state
- Systematic treatment [Lüscher, 1985-1991](#)
 - Conceptually satisfactory basis to study resonances
 - Coupling as parameter (related to width)
- Fit for coupling (assumed constant, related to width)
 - No sensitivity on width (compatible within large error)
 - Small but dominant FV correction for light resonances

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[Go to the continuum](#)

Continuum extrapolation

Goal:

- Eliminate discretization effects

Method:

- Formally in our action: $O(\alpha_s a)$ and $O(a^2)$
 - **But:** discretization effects are tiny
 - Not possible to distinguish between $O(a)$ and $O(a^2)$
- include both in systematic error

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Systematic uncertainties

Goal:

- Accurately estimate total systematic error

Method:

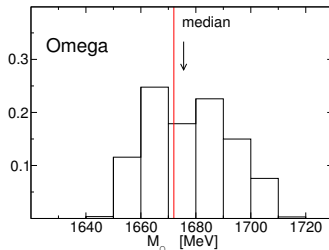
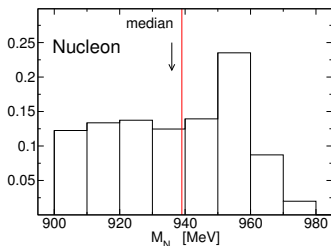
- We account for all the above mentioned effects
- When there are a number of sensible ways to proceed, we take them: Complete analysis for each of
 - 18 fit range combinations
 - ratio/nonratio fits (r_X resp. M_X)
 - $O(a)$ and $O(a^2)$ discretization terms
 - NLO χ PT M_π^3 and Taylor M_π^4 chiral fit
 - 3 χ fit ranges for baryons: $M_\pi < 650/550/450$ MeV

resulting in 432 predictions for each hadron mass with each 2000 bootstrap samples for each Ξ and Ω scale setting

Systematic uncertainties (ctd.)

Method (ctd.):

- Weigh each of the 432 central values by fit quality Q
 - Median of this distribution \rightarrow final result
 - Central 68% \rightarrow systematic error
- Statistical error from bootstrap of the medians



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Particle spectrum

BMW-collaboration, Science **322** (2008), 1224