Dynamical breakdown of Abelian gauge chiral symmetry by strong Yukawa interactions

Petr Beneš

benes@ujf.cas.cz

Nuclear Physics Institute, Řež near Prague, Czech Republic

based on P.B., Tomáš Brauner, Jiří Hošek: hep-ph/0605147 and CERN Yellow Book hep-ph/0608079 The Standard Model:

EW symmetry breaking by means of the Higgs mechanism

Phenomenologically highly succesfull description However:

Difficult to interpret the wrong sign of the Higgs 'mass' squared ($M^2 < 0$)

Fermion masses \propto Yukawa couplings

Our proposal

The EW symmetry breaking can be achieved dynamically:

- Scalars and Yukawa couplings still present in theory
- However no condensing scalars ($M^2 > 0$)
- Fermion masses are nonlinear functions of Yukawa couplings
- In this presentation we
 - do not discuss the full $SU(2) \times U(1)$ theory of EW interactions
 - present merely a simple U(1) toy model

The model

The Lagrangian:

 $\mathcal{L} = \bar{\psi}_{1} i \not{D} \psi_{1} + \bar{\psi}_{2} i \not{D} \psi_{2} + (D_{\mu}\phi)^{\dagger} (D^{\mu}\phi) - \frac{1}{4} F_{\mu\nu} F^{\mu\nu}$ $-M^{2} \phi^{\dagger}\phi - \frac{1}{2} \lambda (\phi^{\dagger}\phi)^{2}$ $+y_{1} (\bar{\psi}_{1L}\psi_{1R}\phi + \bar{\psi}_{1R}\psi_{1L}\phi^{\dagger}) + y_{2} (\bar{\psi}_{2R}\psi_{2L}\phi + \bar{\psi}_{2L}\psi_{2R}\phi^{\dagger})$ $U(1)_{A} \text{ gauge invariance:}$

$$\psi_1 \longrightarrow e^{+i\theta(x)\gamma_5}\psi_1$$

$$\psi_2 \longrightarrow e^{-i\theta(x)\gamma_5}\psi_2$$

$$\phi \longrightarrow e^{-2i\theta(x)}\phi$$

The case of $M^2 < 0$

Standard Higgs mechanism takes place:Scalar develops VEV:

$$\langle \phi \rangle_0 = \sqrt{\frac{-M^2}{\lambda}}$$

The spectrum:

fermions: $m_1 = y_1 \langle \phi \rangle_0$ $m_2 = y_2 \langle \phi \rangle_0$ scalars: $(\phi = \frac{1}{\sqrt{2}}(\phi_1 + i\phi_2))$ $M_1 = \sqrt{2\lambda} \langle \phi \rangle_0$ $(M_2 = 0)$ gauge boson: $M_A = 2\sqrt{2g} \langle \phi \rangle_0$

The case of $M^2 > 0$

Perturbatively, nothing interesting happens: $\ \, | \ \, \langle \phi \rangle_0 = 0$ Fermions remain massless No mass splitting in the scalar sector $\blacksquare \Rightarrow$ the $U(1)_A$ symmetry remains unbroken But what about some *nonperturbative* treatment? Make use of the Schwinger-Dyson equations Consider only symmetry-breaking 2-point Green functions:

 $\langle \phi \phi \rangle, \langle \phi^{\dagger} \phi^{\dagger} \rangle, \langle \psi_{iL} \overline{\psi}_{iR} \rangle, \langle \psi_{iR} \overline{\psi}_{iL} \rangle$

Schwinger-Dyson equations



– 1PI propagator
– full propagator

Schwinger-Dyson equations

$$\begin{split} \Sigma_{1,p} &= \mathrm{i} y_1^2 \int \frac{\mathrm{d}^4 k}{(2\pi)^4} \frac{\Sigma_{1,k}}{k^2 - \Sigma_{1,k}^2} \frac{\Pi_{k-p}}{[(k-p)^2 - M^2]^2 - |\Pi_{k-p}|^2} \\ \Sigma_{2,p} &= \mathrm{i} y_2^2 \int \frac{\mathrm{d}^4 k}{(2\pi)^4} \frac{\Sigma_{2,k}}{k^2 - \Sigma_{2,k}^2} \frac{\Pi_{k-p}^*}{[(k-p)^2 - M^2]^2 - |\Pi_{k-p}|^2} \\ \Pi_p &= -\sum_{j=1,2} 2\mathrm{i} y_j^2 \int \frac{\mathrm{d}^4 k}{(2\pi)^4} \frac{\Sigma_{j,k}}{k^2 - \Sigma_{j,k}^2} \frac{\Sigma_{j,k-p}}{(k-p)^2 - \Sigma_{j,k-p}^2} \\ &+ \mathrm{i} \lambda \int \frac{\mathrm{d}^4 k}{(2\pi)^4} \frac{\Pi_k}{(k^2 - M^2)^2 - |\Pi_k|^2} \end{split}$$

Notation: $\Sigma_{1,p} \equiv \Sigma_1(p^2)$, etc.

The spectrum

Fermions:
$$m_1^2 = \Sigma_1^2(m_1^2)$$

 $m_2^2 = \Sigma_2^2(m_2^2)$
Scalars: $M_1^2 = M^2 + |\Pi(M_1^2)$
 $M_2^2 = M^2 - |\Pi(M_2^2)$

Gauge boson:

- Need to employ Ward identities
- 'Would-be' Goldstone boson is composite of fermions and scalars

• $M_A^2 = g^2 \times$ complicated integral of Σ_1, Σ_2, Π

Unlike the Higgs case the masses are related through dependence on Σ_1, Σ_2, Π

Numerical computation

SD equations still too difficult, need some simplifications:

- Switch to Euclidean metric
 some poles removed (but not all)
- Set $\lambda = 0$

Then:

- For some values (y₁, y₂) the solution to the SD equations was found
- Self-energies Σ_1, Σ_2, Π are rapidly decreasing functions

The *y*_{1,2}-dependence of the spectrum was probed
The case of fermions is most interesting . . .

Fermionic spectrum



Calculated for fixed $y_2 = 88$

We have found that

- the Model is capable of providing SSB
- the ratio of dynamically generated masses can be large for not vastly different Yukawa couplings

But still much to do:

- Get some analytical insight into the above results
- Try to find solution to the SD equations in the Minkowski space
- Work out the realistic SU(2) × U(1) model (sketch already present in hep-ph/0407339)