



MeV-scale Reheating and and Neutrino Thermalization

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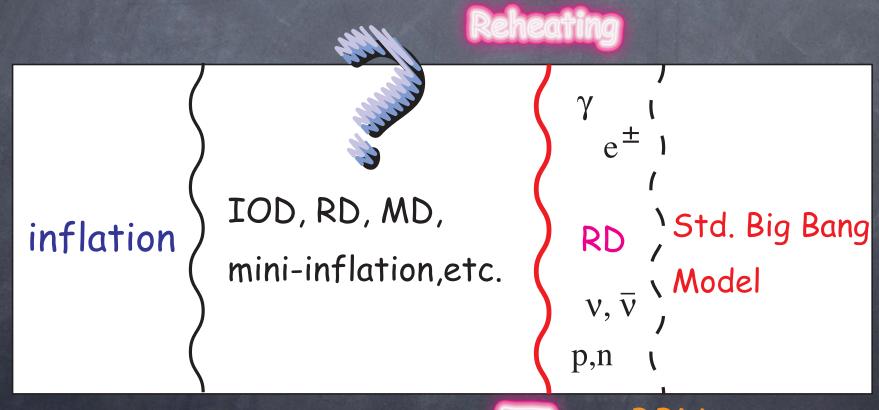
(DESY, Theory Group & Univ. of Tokyo, ICRR)

collaborators: K. Ichikawa and M. Kawasaki

K. Ichikawa, M. Kawasaki, F.T., Phys.Rev. D72(2005)043522

I. Introduction

What do we know about thermal history of the universe?





BBN

Big bang nucleosynthesis

n-p transformation decouples when

$$\left(\Gamma_{n \leftrightarrow p} = H \right)$$

$$n \leftrightarrow p + e^{-} + \bar{\nu}_{e}$$

$$n + \nu_{e} \leftrightarrow p + e^{-}$$

$$n + e^{+} \leftrightarrow p + \bar{\nu}_{e}$$

$$\left(\frac{n}{p}\right)_{\text{EQ}} = \exp\left(-\frac{Q}{T}\right)$$

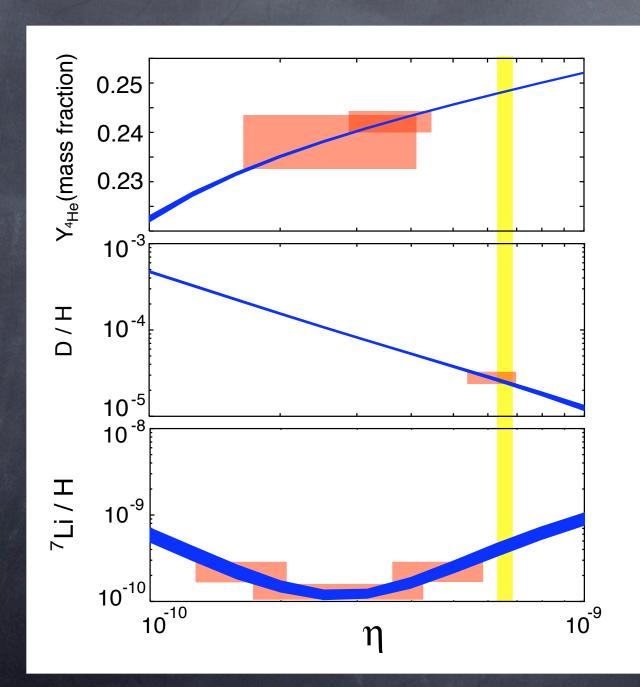
$$Q = m_n - m_p = 1.293 \text{MeV}$$
$$T \sim 0.8 \text{MeV}$$

$$2n + 2p \rightarrow 4 \text{He}$$

Almost all neutrons are absorbed in ⁴He!

$$B(^{4}\text{He}) = 28.3 \text{MeV} \gg B(D, T, ^{3}\text{He})$$

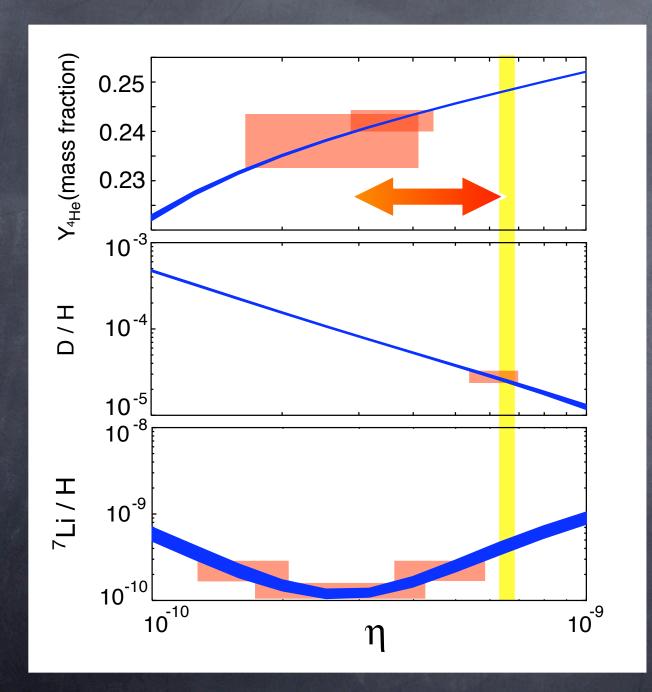
Present status of BBN (th.&obs.)



sBBN includes only one free parameter:

$$\eta = rac{n_B}{n_{m{\gamma}}}$$

Present status of BBN (th.&obs.)



- To cure the discrepancy, we need
 - large L
 or

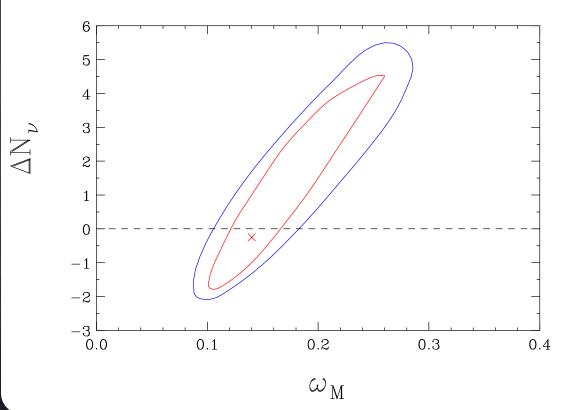
$$\boxed{N_{\nu} < 3}$$

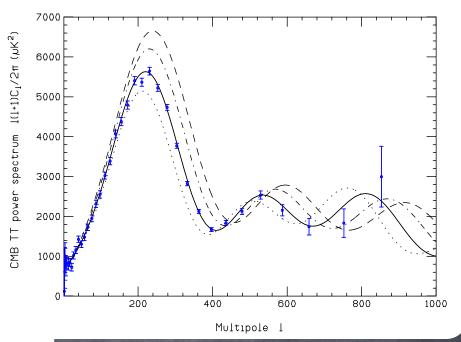
realized if
TR = O(MeV)

 $N_{\nu} \equiv \Sigma \rho_{\nu} / \rho_{\nu, \text{std.}}$

Cosmic Microwave Background

The form of Cl also depends on ΔN .





 $0.9 \le N \le 8.3$ Nbest fit = 2.75

from WMAP data

V. Barger et al, PLB 566(2003)8



What if TRH is around MeV?

Neutrinos might not be fully thermalized. [Kawasaki, Kohri, Sugiyama, '99, '00]



The predicted abundances of the light elements (especially 4He) are affected.

It has been widely believed that $N_
u$ Y_p





What we did:

We investigated the MeV-scale reheating scenario, taking account of

thermalization processes of neutrinos

(+ neutrino oscillations)

 \checkmark We found that $^4{\rm He}$ abundance (increases) as N_{ν} decreases due to flavor mixings!



(to be detailed later)

II. Role of neutrinos in BBN

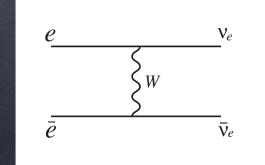
Reheating processes (Setup)

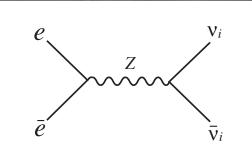
$$\phi \rightarrow \gamma, e^{\pm}$$

We define the reheating temperature by $\Gamma=3H(T_R)$

$$\Gamma = 2.03 \left(\frac{T_R}{\text{MeV}}\right)^2 \text{ sec}^{-1}.$$

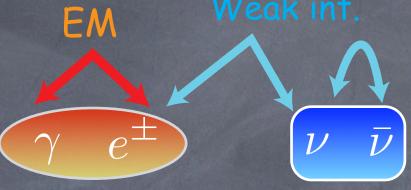
• We assume that neutrinos are exclusively produced via $e^- + e^+ \leftrightarrow \nu + \bar{\nu}$





Decoupling of u

radiation:



thermal. eq. via EM int.

decouple @ T = a few MeV

 $T < m_{\mu}$

In std. BBC,

$$T_d \sim 5 {
m MeV}$$
 for $u_{\mu, au}$ $T_d \sim 3 {
m MeV}$ for u_e

[Hannestad & Madsen '95] [Dolgov, Hansen & Semikoz, '97, '99]

 ν_e is more likely to be produced at T ~MeV.

\bullet What if N_{ν} < 3 ?

$$N_{\nu} \equiv \Sigma \rho_{\nu} / \rho_{\nu, \text{std.}}$$

1. Hubble parameter decreases.



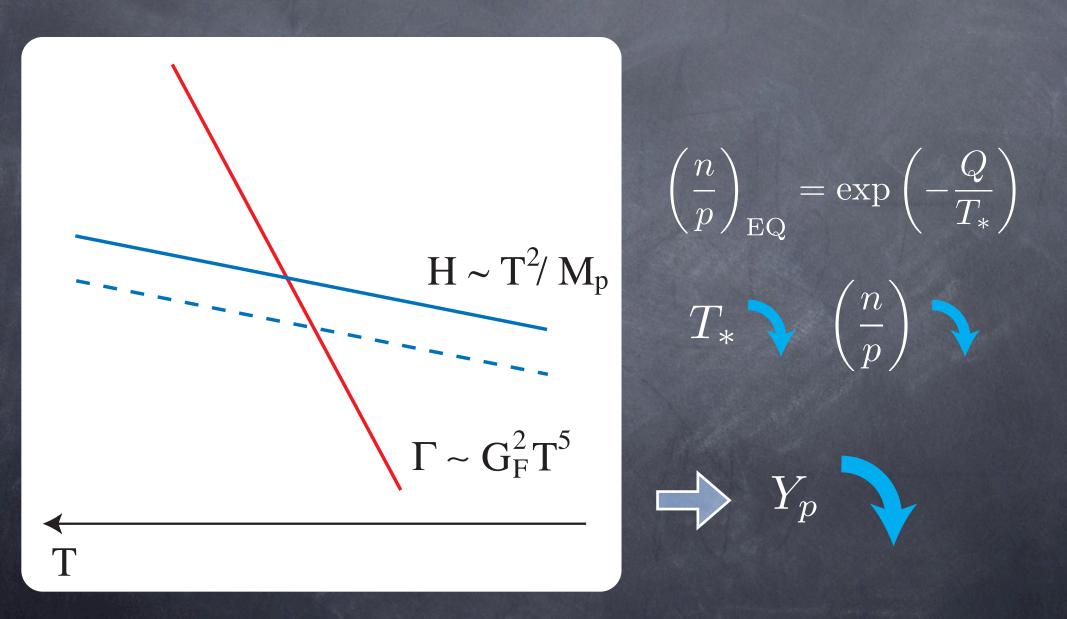
2. The n-p tranf. rate also decreases.

$$\Gamma_{n\leftrightarrow p}$$
 Y_p

3. D bottleneck opens later.

$$H(T_D)^{-1}$$
 Y_p

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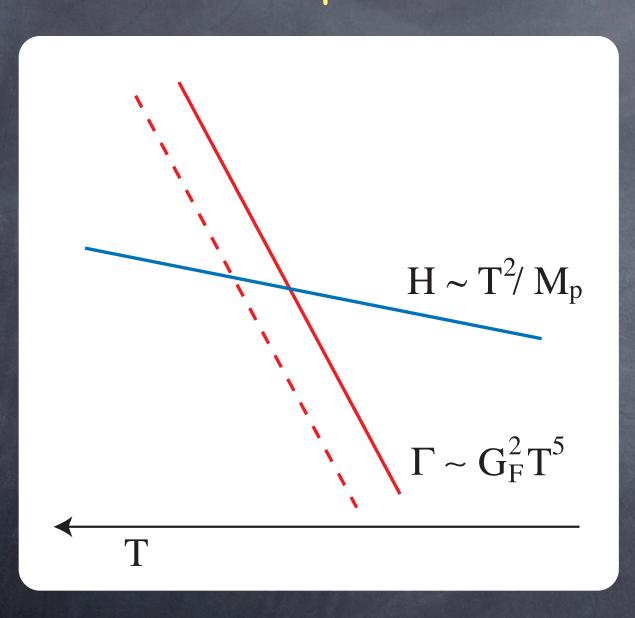
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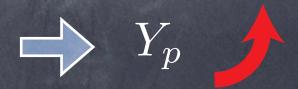


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$$\left(\frac{n}{p}\right)_{\text{EQ}} = \exp\left(-\frac{Q}{T_*}\right)$$



\bullet What if N_{ν} < 3 ?

$$N_{\nu} \equiv \Sigma \rho_{\nu} / \rho_{\nu, \text{std.}}$$

✓ Hubble parameter decreases.







✓ The n-p tranf. rate also decreases.

$$\Gamma_{n\leftrightarrow p}$$
 Y_p





D bottleneck opens later.

$$H(T_D)^{-1}$$



III. Thermalization of Neutrinos

Density matrices:

$$\left\langle a_j^{\dagger}(\mathbf{p}) a_i(\mathbf{p}') \right\rangle = (2\pi)^3 \delta^{(3)}(\mathbf{p} - \mathbf{p}') \left[\rho_{\mathbf{p}} \right]_{ij},$$

$$\left\langle b_i^{\dagger}(\mathbf{p}) b_j(\mathbf{p}') \right\rangle = (2\pi)^3 \delta^{(3)}(\mathbf{p} - \mathbf{p}') \left[\overline{\rho}_{\mathbf{p}} \right]_{ij}.$$

$$\{i,j\} = \{e,\mu,\tau\}$$

- \circ diagonal components: dist. func. of u_i
- \circ off-diagonal ones: corr. between u_i and u_j .

QKEs for density matrix

$$\nu + e^{\pm} \leftrightarrow \nu + e^{\pm}$$
$$\nu + \bar{\nu} \leftrightarrow e^{-} + e^{+}$$

$$Ha rac{d
ho_p}{da} = -i[\Omega(p),
ho_p] + I_{
m coll}(p).$$
 refractive term [flavor mixings] collision term [neutrino there

[flavor mixings] [neutrino thermalization]

where

$$\Omega(p) \equiv \Omega_V(p) - \frac{8\sqrt{2G_F p}}{3m_W^2} E$$

$$\Omega_V(p) = rac{1}{2p} U M^2 U^T, \qquad M^2 \equiv \left(egin{array}{cc} m_1^2 & 0 \ 0 & m_2^2 \end{array}
ight), \qquad U \equiv \left(egin{array}{cc} \cos heta_{12} & \sin heta_{12} \ -\sin heta_{12} & \cos heta_{12} \end{array}
ight).$$

$$E = \left(\begin{array}{cc} \rho_e + \rho_{\bar{e}} & 0 \\ 0 & 0 \end{array}\right) = \left(\begin{array}{cc} (7/60)\pi^2 T^4 & 0 \\ 0 & 0 \end{array}\right), \left(\begin{array}{c} \text{Approximate 2 flavor analysis assuming } \theta_{13} \approx 0 \end{array}\right)$$

Parameters:

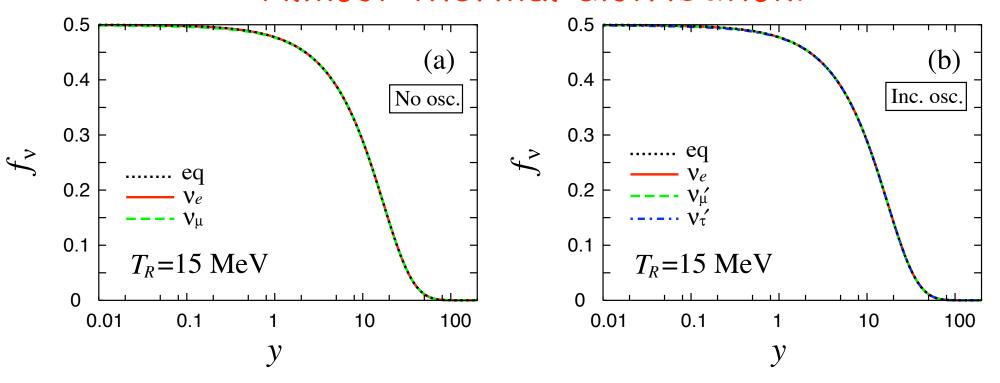
$$\sin^2 \theta_{12} = 0.315$$

$$m_2^2 - m_1^2 = 7.3 \times 10^{-5} \text{eV}^2$$

IV. Results & Discussions

The final dist. functions [T_R=15MeV]





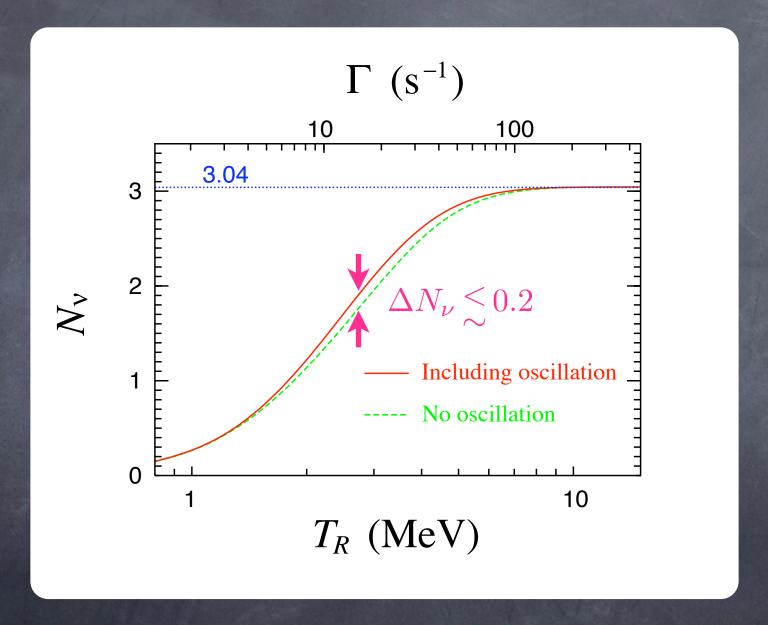
The final dist. functions [T_R=2.5MeV]

$$f_{\nu_e} > f_{\nu'_{\mu}} = f_{\nu'_{\tau}}$$

$$f_{\nu_e} \sim f_{\nu'_\mu} > f_{\nu'_\tau}$$

Due to flavor mixings, $f_{\nu'_{\mu}} - f_{\nu_{e}} - \Gamma_{n \leftrightarrow p}$

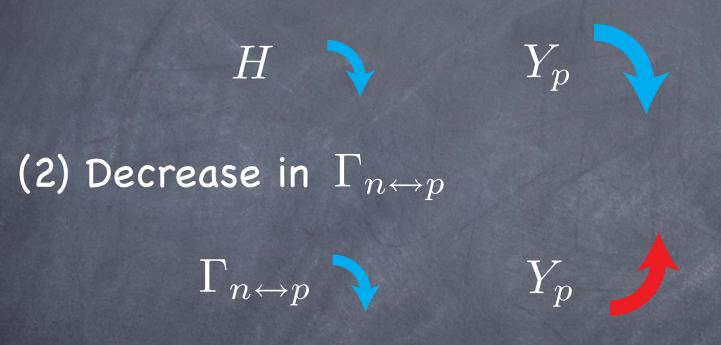
Effective number of neutrinos



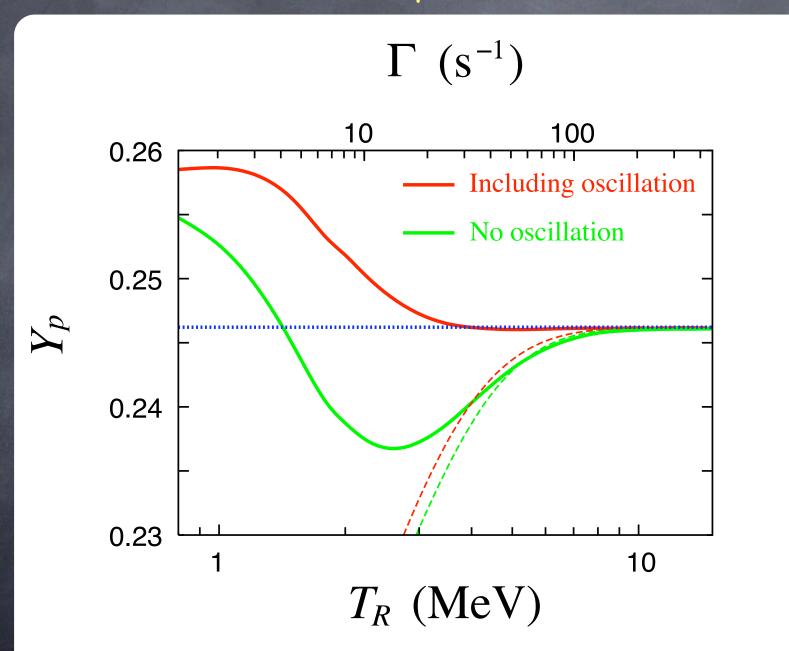
 N_{ν} (i.e. H) does not change much!

Two competing effects:

(1) Slowing down of the expansion rate



Taking into consideration the neutrino oscillations, the effect (2) becomes more important!



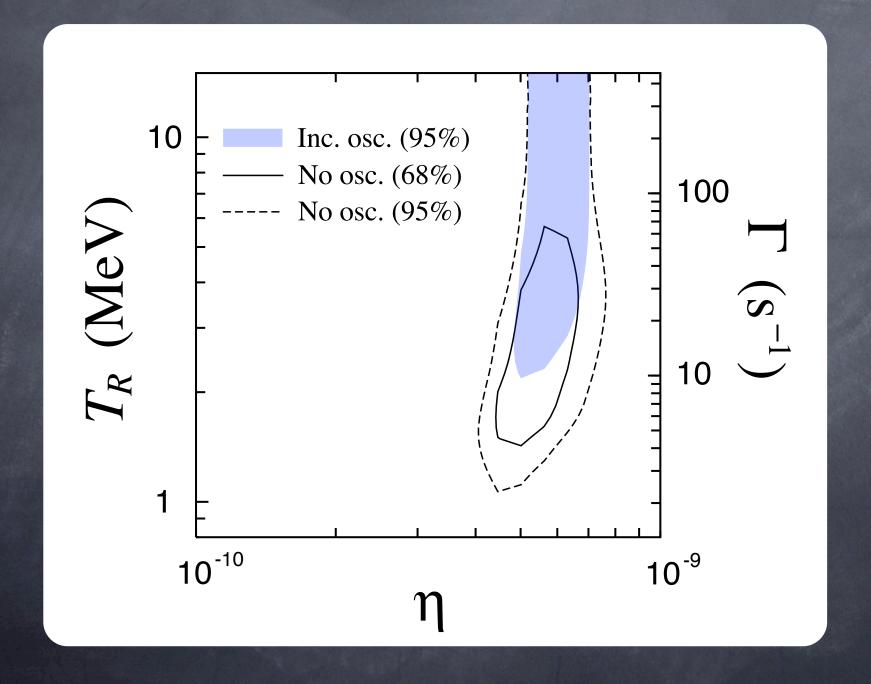
V. Conclusion

- We have investigated the MeV-scale reheating scenario, paying particular attention to neutrino oscillations.
- In contrast to the widespread picture,

$$N_{
u}$$
 Y_p

while T_R decreases.

 σ $T_R \gtrsim 2\,\mathrm{MeV}$ and $N_
u \gtrsim 1.2\,\mathrm{was}$ obtained.



Important reactions in BBN

